

First-principles modeling of efficiency of halide perovskites

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Van de Walle Computational Materials Group

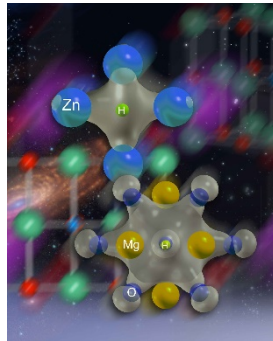
vandewalle.materials.ucsb.edu

First-principles calculations

Density functional theory, many-body perturbation theory

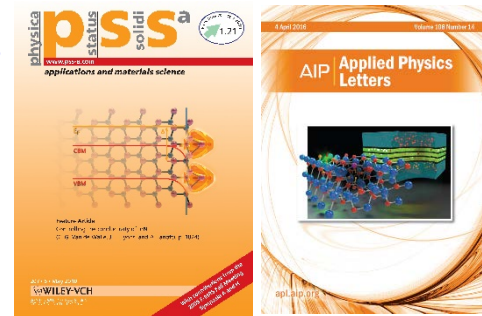
Oxides

- Transparent conductors
- Dielectrics
- Thermal barriers
- Complex oxides
- Power electronics



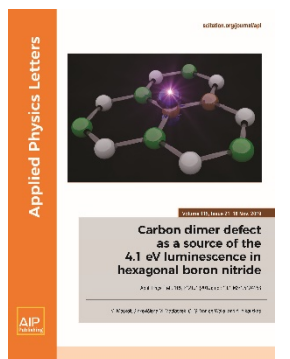
Nitrides

- Doping
- Surfaces
- Interfaces
- Efficiency, loss



Quantum defects

- Qubits
- Single photon emitters
- Decoherence



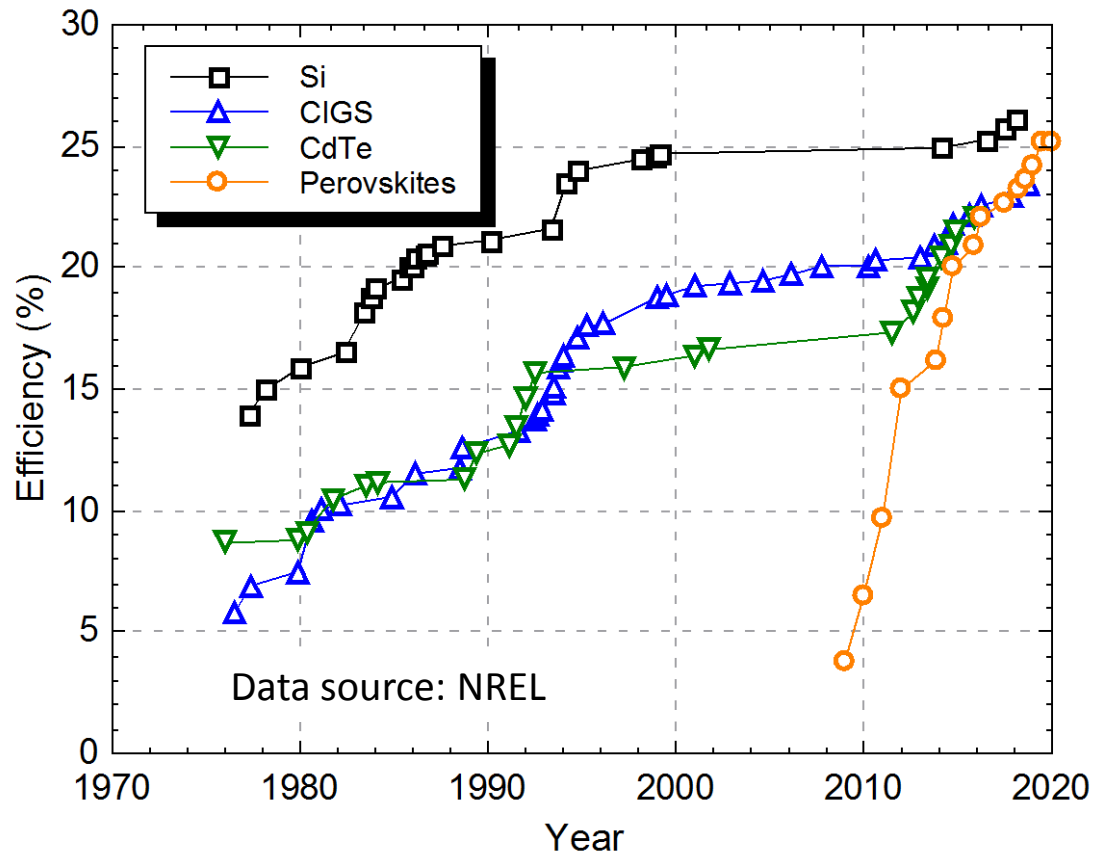
Hybrid perovskites

- Recombination mechanisms
- Defects
- Impurities
- Efficiency

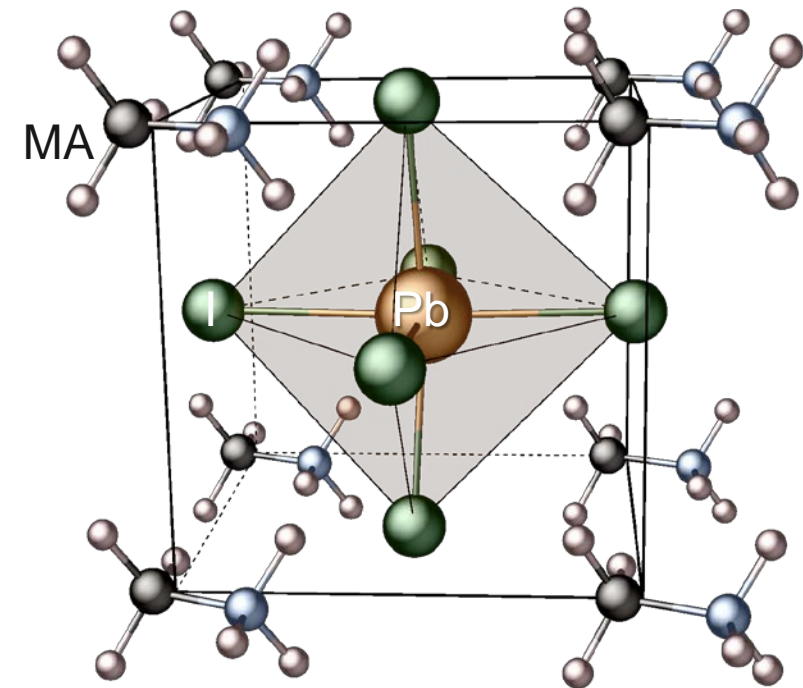


Halide perovskites

- Efficient optoelectronic materials
 - Solar cells, light emitting diodes (LEDs)

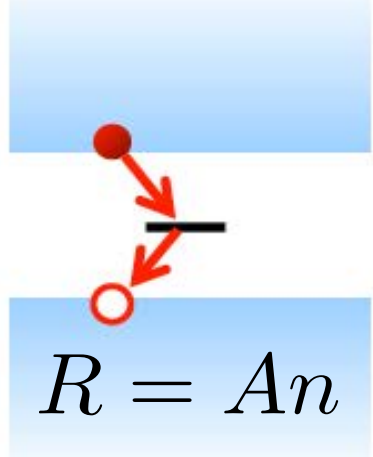


- Prototype: methylammonium lead iodide
 MAPbI_3
 - A: Cs^+ , MA^+ , FA^+
 - B: Pb^{2+} , Sn^{2+}
 - X: I^- , Br^- , $\text{Cl}^- \dots$
- General: ABX_3



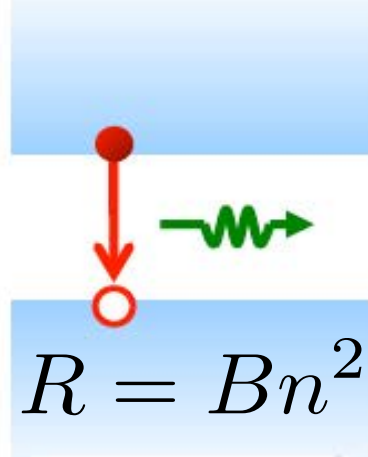
Recombination mechanisms

Defect-assisted



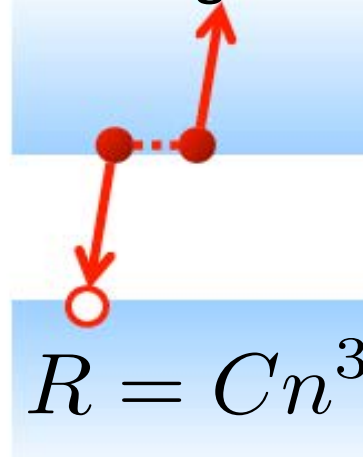
$$R = An$$

Radiative



$$R = Bn^2$$

Auger

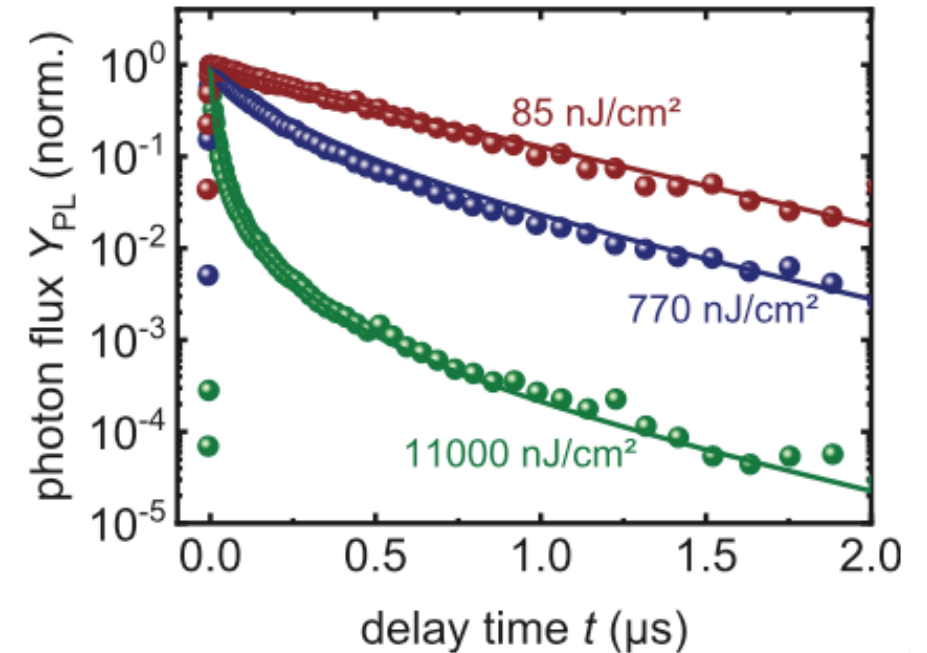


$$R = Cn^3$$

“Shockley-Read-Hall”
(SRH)

n : carrier density

Rate equation:
$$\frac{dn}{dt} = -An - Bn^2 - Cn^3$$

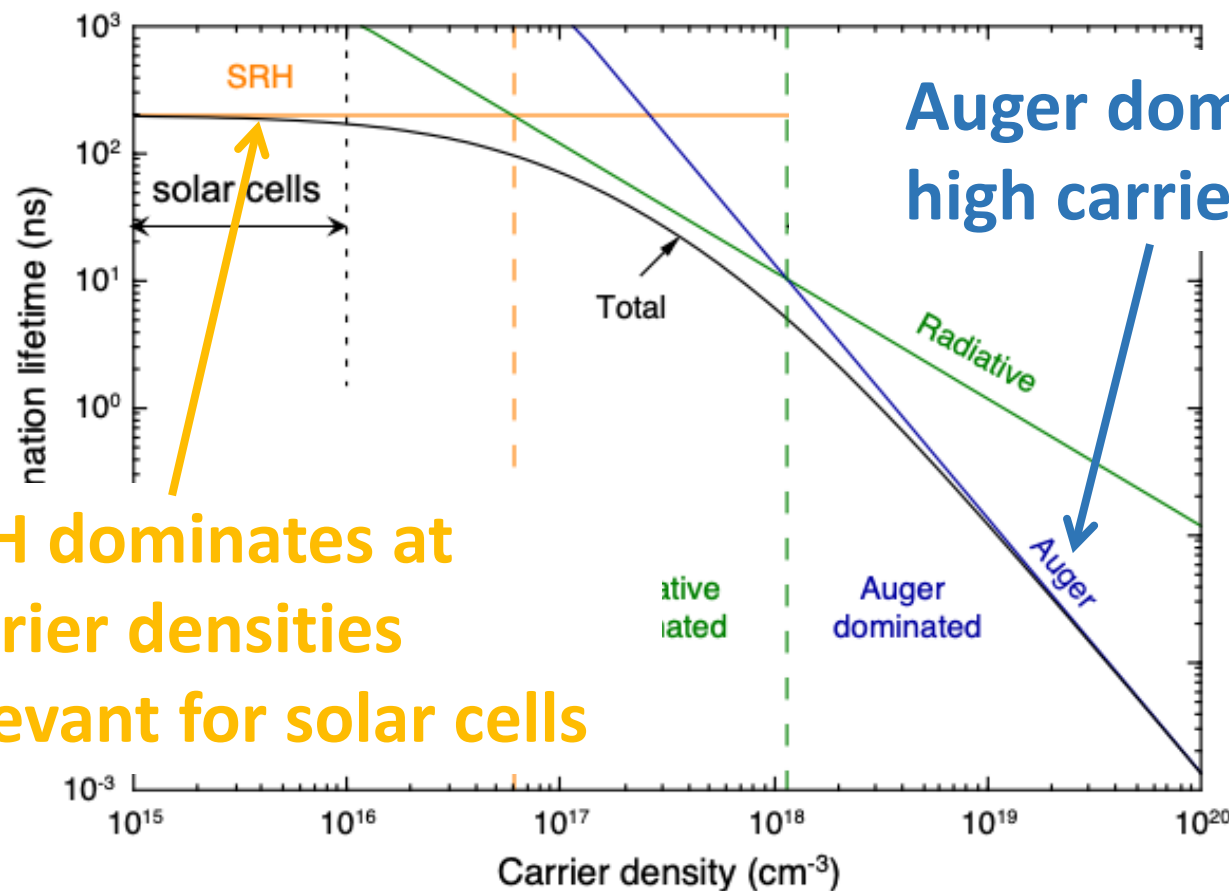


T. Kirchartz *et al.*,
Adv. Energy Mater. **10**, 1904134 (2020).

Fitting to experimental data can introduce uncertainties, and does not elucidate fundamental mechanisms

Recombination in halide perovskites

First-principles studies of recombination rates



SRH dominates at carrier densities relevant for solar cells

Auger dominates at high carrier densities

X. Zhang *et al.*, J. Phys. Chem. Lett. **9**, 2903 (2018).

X. Zhang *et al.*, ACS Energy Lett. **3**, 2329 (2018).

J.-X. Shen *et al.*, Adv. Energy Mater. **8**, 1801027 (2018).

X. Zhang *et al.*, Adv. Energy Mater. **10**, 1902830 (2020).

Focus of this talk

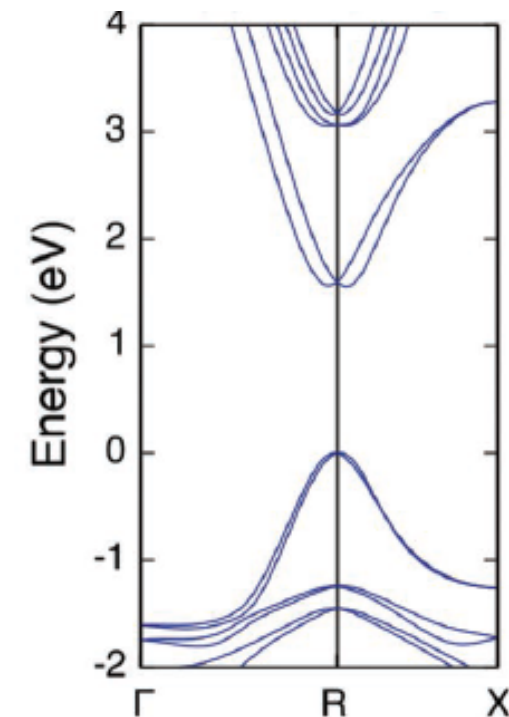
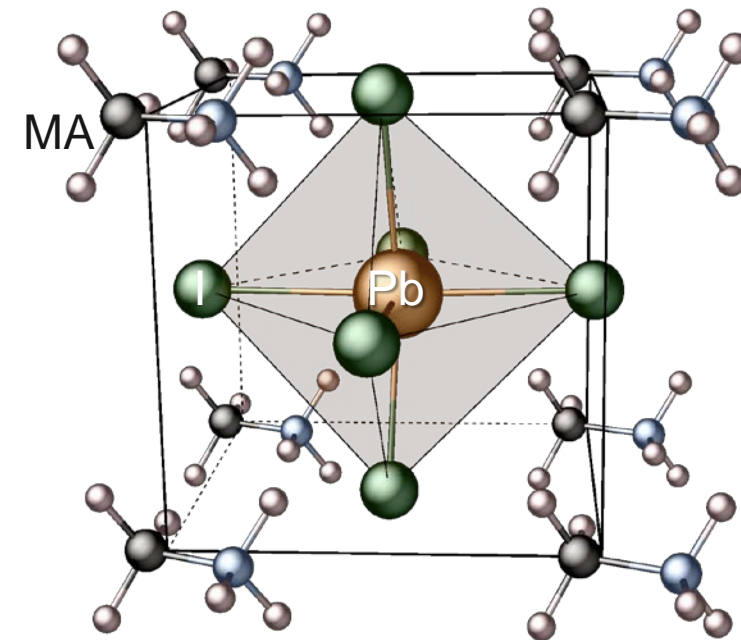
- First-principles studies
 - Radiative recombination
 - Auger recombination
 - Defect-assisted Shockley-Read-Hall (SRH) recombination

- First-principles approach:
 - Density functional theory

- HSE hybrid functional

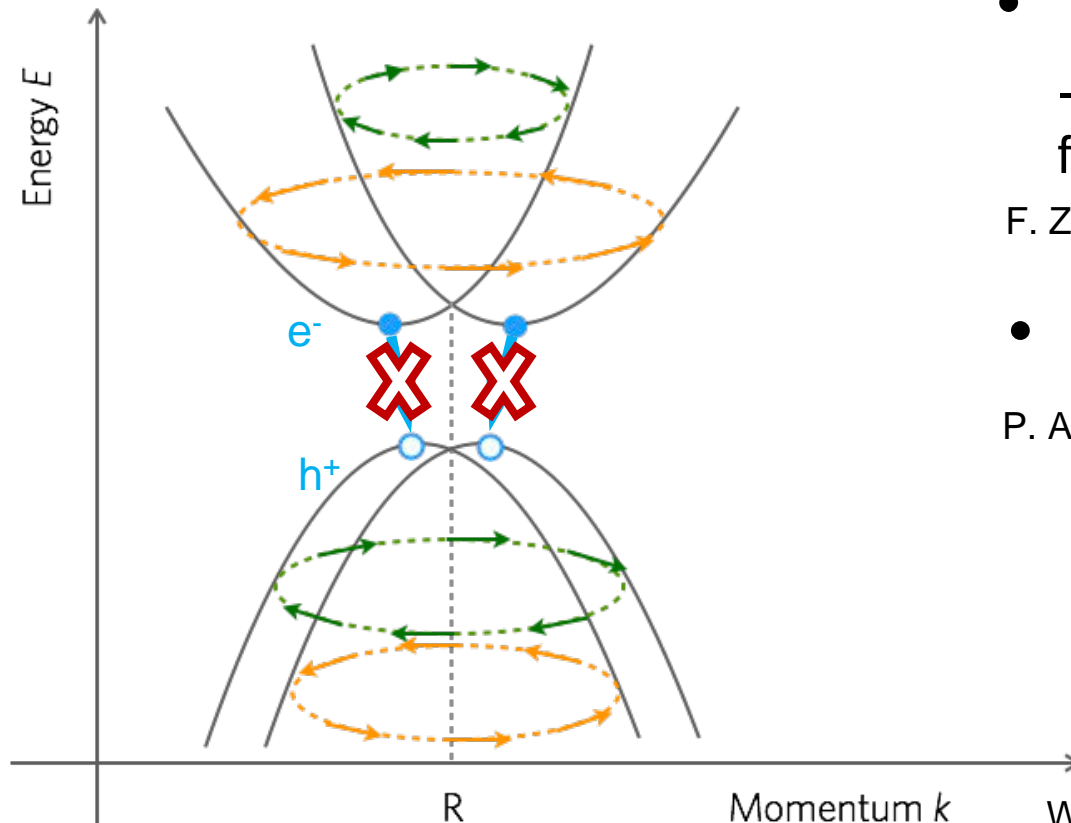
J. Heyd *et al.*, J. Chem. Phys. **118**, 8207 (2003).

- Vienna Ab-initio Simulation Package (VASP);
Quantum Espresso



Slow radiative recombination?

Rashba spin-orbit coupling



- Spin mismatch
 - Phenomenological model with fitted parameters

F. Zheng *et al.*, Nano Lett. **15**, 7794 (2015).

- Momentum mismatch

P. Azarhoosh *et al.*, APL Mater. **4**, 091501 (2016).

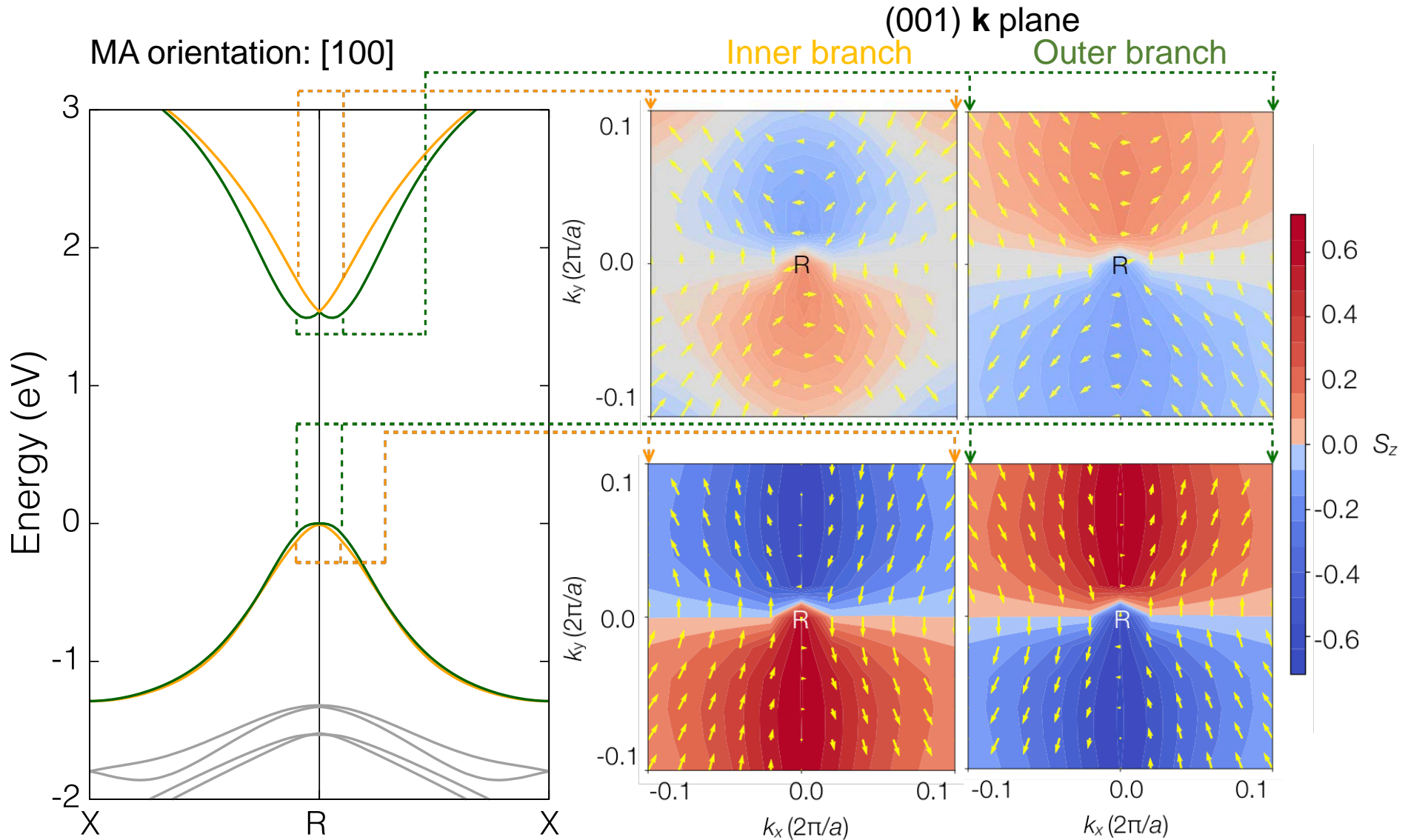
MAPbI₃: $\sim 10^{-13} \text{ cm}^3\text{s}^{-1}$

GaAs: $\sim 10^{-9} \text{ cm}^3\text{s}^{-1}$

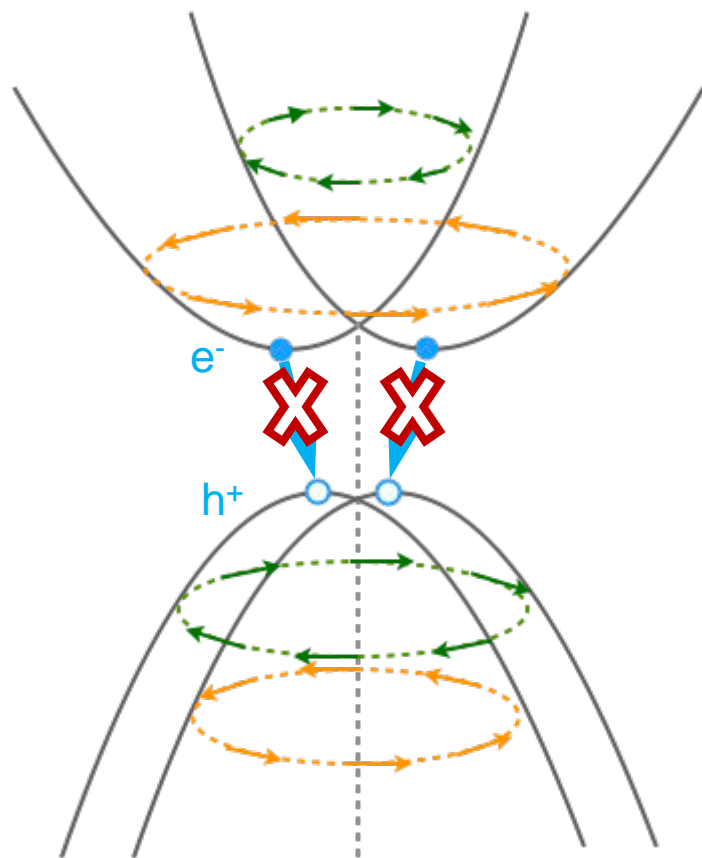
Si: $\sim 10^{-14} \text{ cm}^3\text{s}^{-1}$

W. Tress, Adv. Energy Mater. **7**, 1602358 (2017).

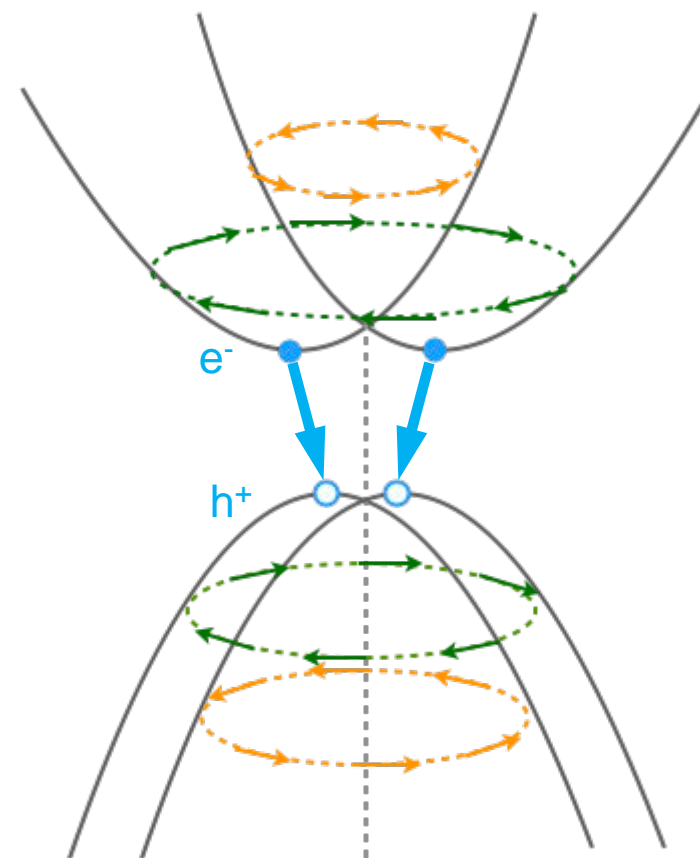
First-principles spin texture



Phenomenological
spin orientation



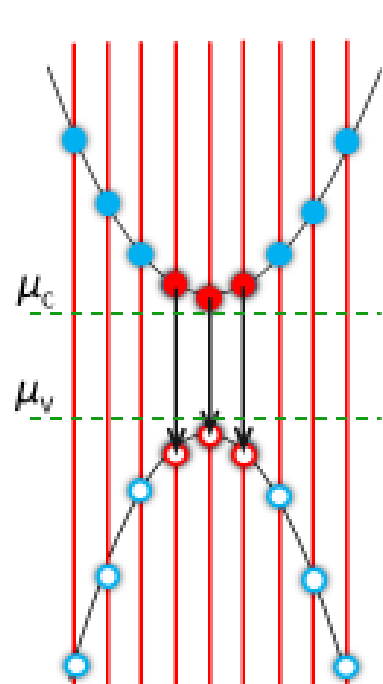
vs. First principles
spin orientation



Radiative recombination

➤ Fermi's Golden Rule:

$$B = \frac{n_r e^2}{\pi \epsilon_0 m_e^2 c^3 \hbar^2 n^2 V} \sum_{c\mathbf{k}} \underbrace{f_{c\mathbf{k}} (1 - f_{v\mathbf{k}})}_{\text{Quasi-Fermi occupation}} \underbrace{(\epsilon_{c\mathbf{k}} - \epsilon_{v\mathbf{k}})}_{\text{Energy difference}} \underbrace{|\mathbf{M}_{cv\mathbf{k}}|^2}_{\text{Matrix elements}}$$



Prefactor

Energy difference

$$f_{c\mathbf{k}} = \frac{1}{1 + e^{\frac{E_{c\mathbf{k}} - \mu_c}{k_B T}}}$$

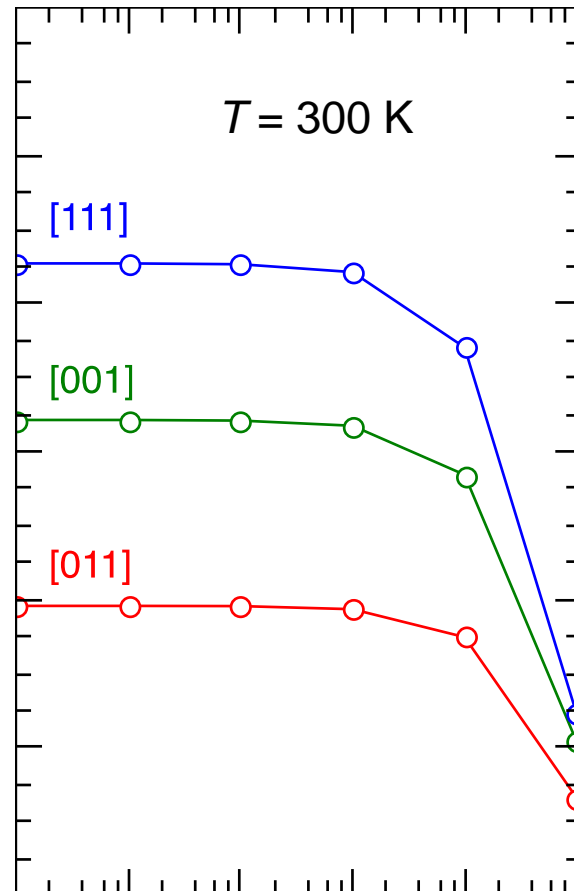
$$f_{v\mathbf{k}} = \frac{1}{1 + e^{\frac{E_{v\mathbf{k}} - \mu_v}{k_B T}}}$$

$$|\mathbf{M}_{cv\mathbf{k}}|^2 = \frac{1}{3} \sum_i |\langle \psi_c | p_i | \psi_v \rangle|^2$$

$i = x, y, z$

$$\sum_{c\mathbf{k}} f_{c\mathbf{k}} = \sum_{v\mathbf{k}} (1 - f_{v\mathbf{k}}) = n N_{\mathbf{k}} V_{\text{cell}}$$

Radiative recombination



- Weak dependence on MA orientation (factor of 2)
- Limited impact of momentum mismatch on radiative recombination
- High radiative recombination coefficients ($\sim 10^{-10}\text{ cm}^3\text{s}^{-1}$)
- Promising for light-emitting applications

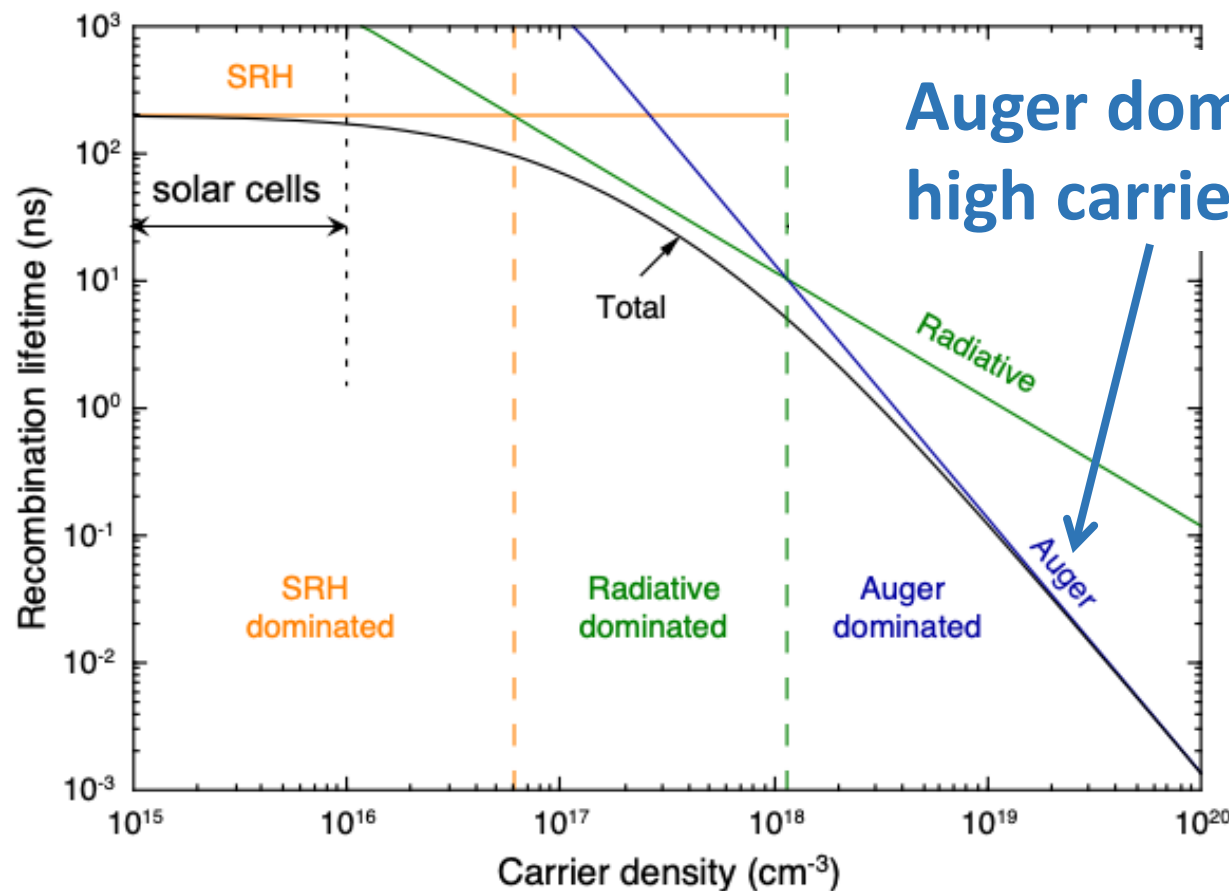
H. Cho *et al.*, *Science* **350**, 1222 (2015).

Y.-H. Kim *et al.*, *Adv. Mater.* **27**, 1248 (2015).

X. Zhang, J.-X. Shen, W. Wang, and C. G. Van de Walle, *ACS Energy Lett.* **3**, 2329 (2018).

Recombination in halide perovskites

First-principles studies of recombination rates



X. Zhang *et al.*, J. Phys. Chem. Lett. **9**, 2903 (2018).

X. Zhang *et al.*, ACS Energy Lett. **3**, 2329 (2018).

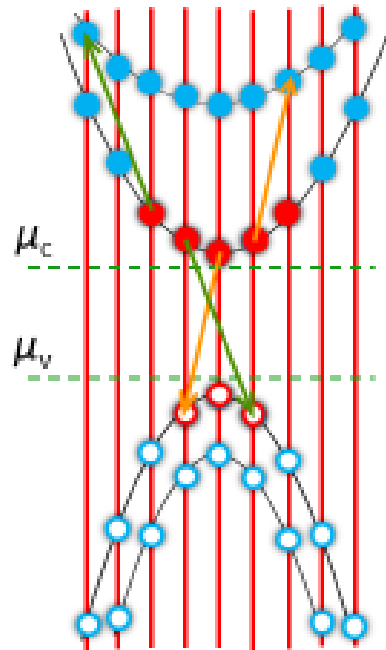
J.-X. Shen *et al.*, Adv. Energy Mater. **8**, 1801027 (2018).

X. Zhang *et al.*, Adv. Energy Mater. **10**, 1902830 (2020).

Auger recombination

➤ Fermi's Golden Rule:

$$C = \frac{2\pi}{\hbar n^3} \sum_{1234} \underbrace{f_1 f_2 (1 - f_3) (1 - f_4)}_{\text{Quasi-Fermi occupation}} \underbrace{|M_{1234}|^2}_{\text{Matrix elements}} \underbrace{\delta(\varepsilon_1 + \varepsilon_2 - \varepsilon_3 - \varepsilon_4)}_{\text{Energy conservation}}$$



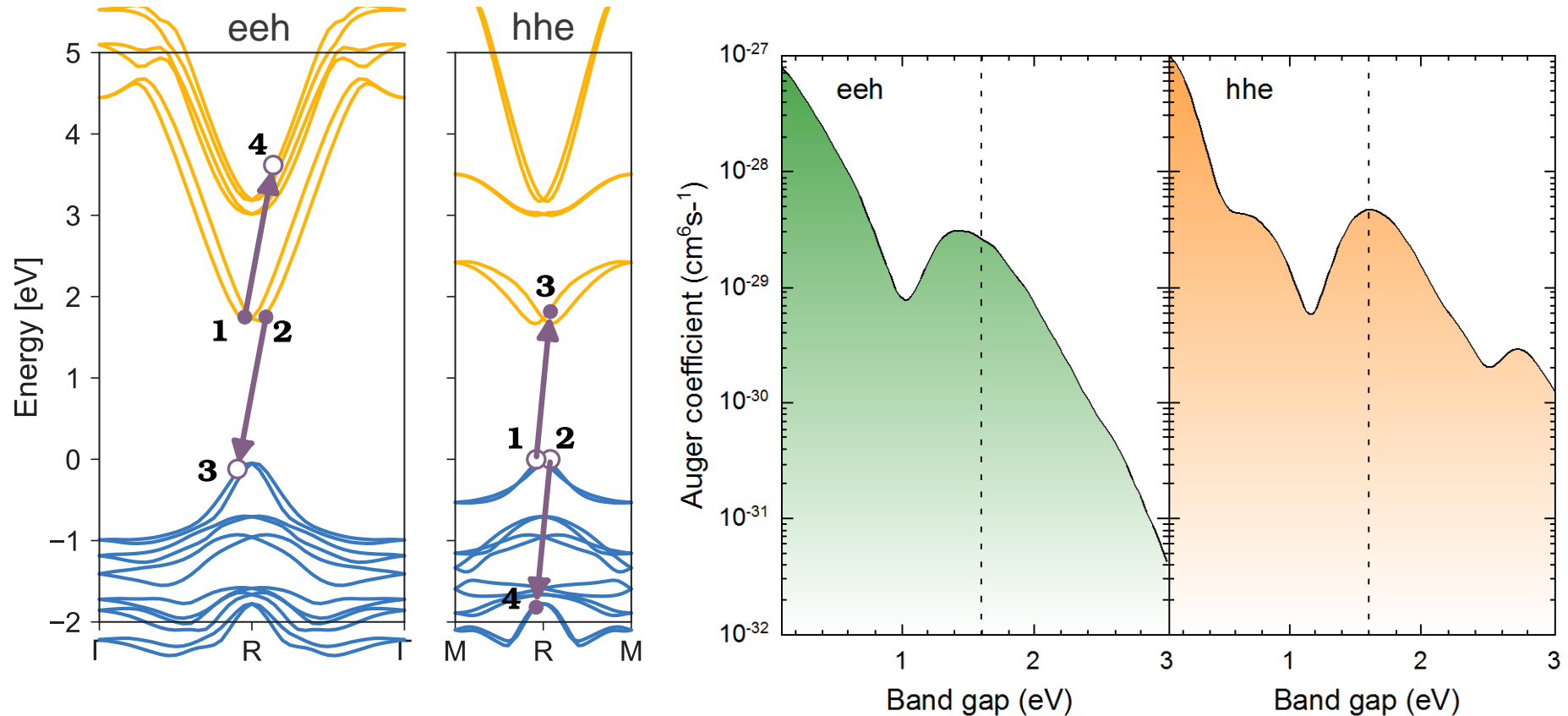
- First-principles calculations for very dense **k**-point grid, e.g., 50 x 50 x 50 for sampling the first Brillouin zone
- Directly search for all possible Auger events that conserve energy and momentum
- Usually on the order of a few tens of millions of possible events

$$|M_{1234}|^2 = |M_{1234}^d - M_{124}^x|^2 + |M_{124}^d|^2 + |M_{1234}^x|^2$$

Direct process: $M_{1234}^d = \langle \psi_1 \psi_2 | W | \psi_3 \psi_4 \rangle$

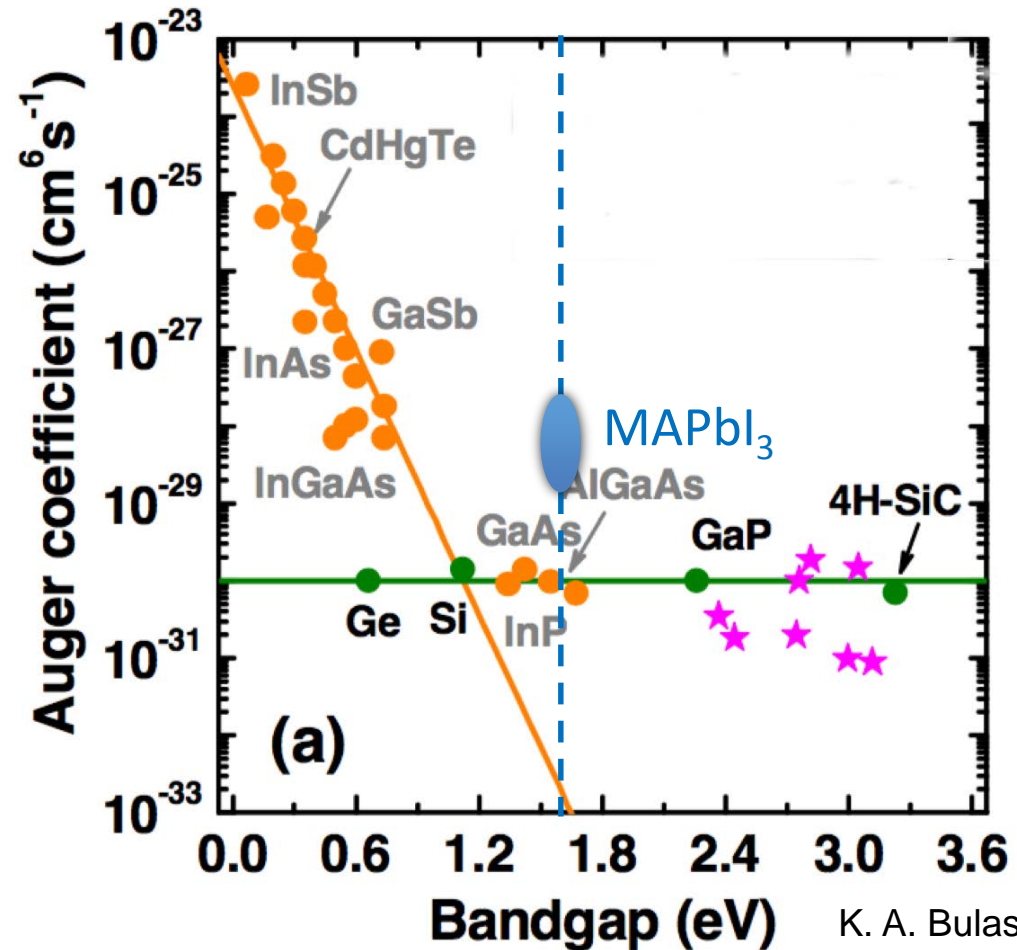
Exchange process: $M_{1234}^x = \langle \psi_1 \psi_2 | W | \psi_4 \psi_3 \rangle$

Auger recombination



- Theory: $C_{\text{tot}} = 7 \times 10^{-29} \text{ cm}^6\text{s}^{-1}$ at $E_g = 1.6 \text{ eV}$
- Exp.: $10^{-29} \sim 10^{-28} \text{ cm}^6\text{s}^{-1}$ [R. L. Milot *et al.*, *Adv. Funct. Mater.* **25**, 6218 (2015).]

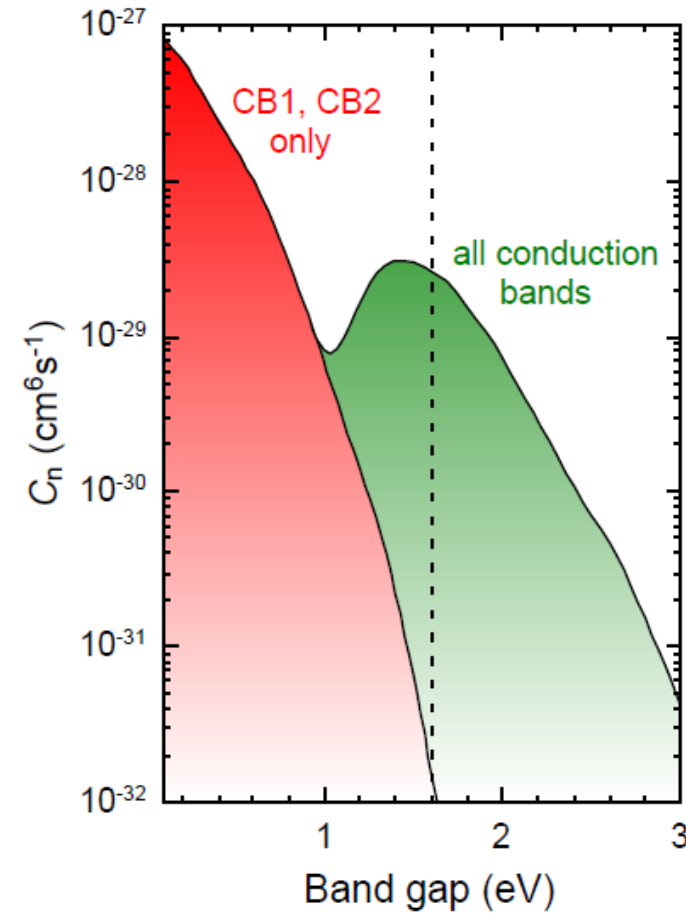
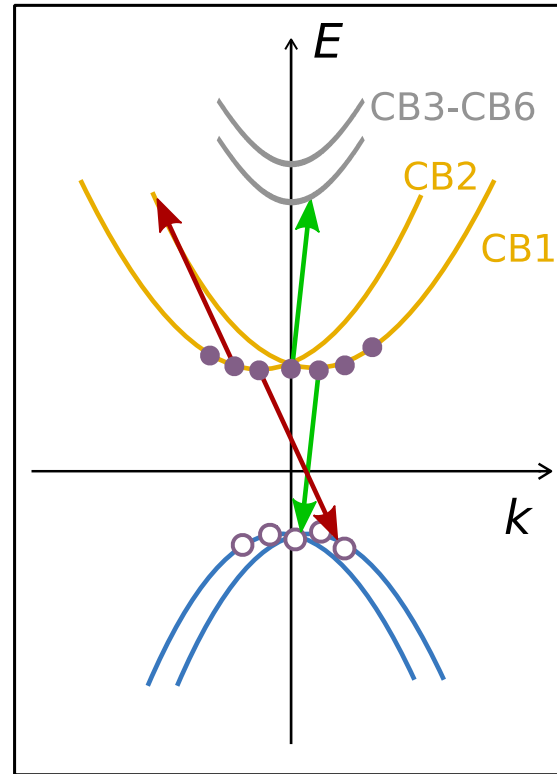
Auger recombination



K. A. Bulashevich *et al.*, Phys. Stat. Sol. (c) **5**, 2066 (2008).

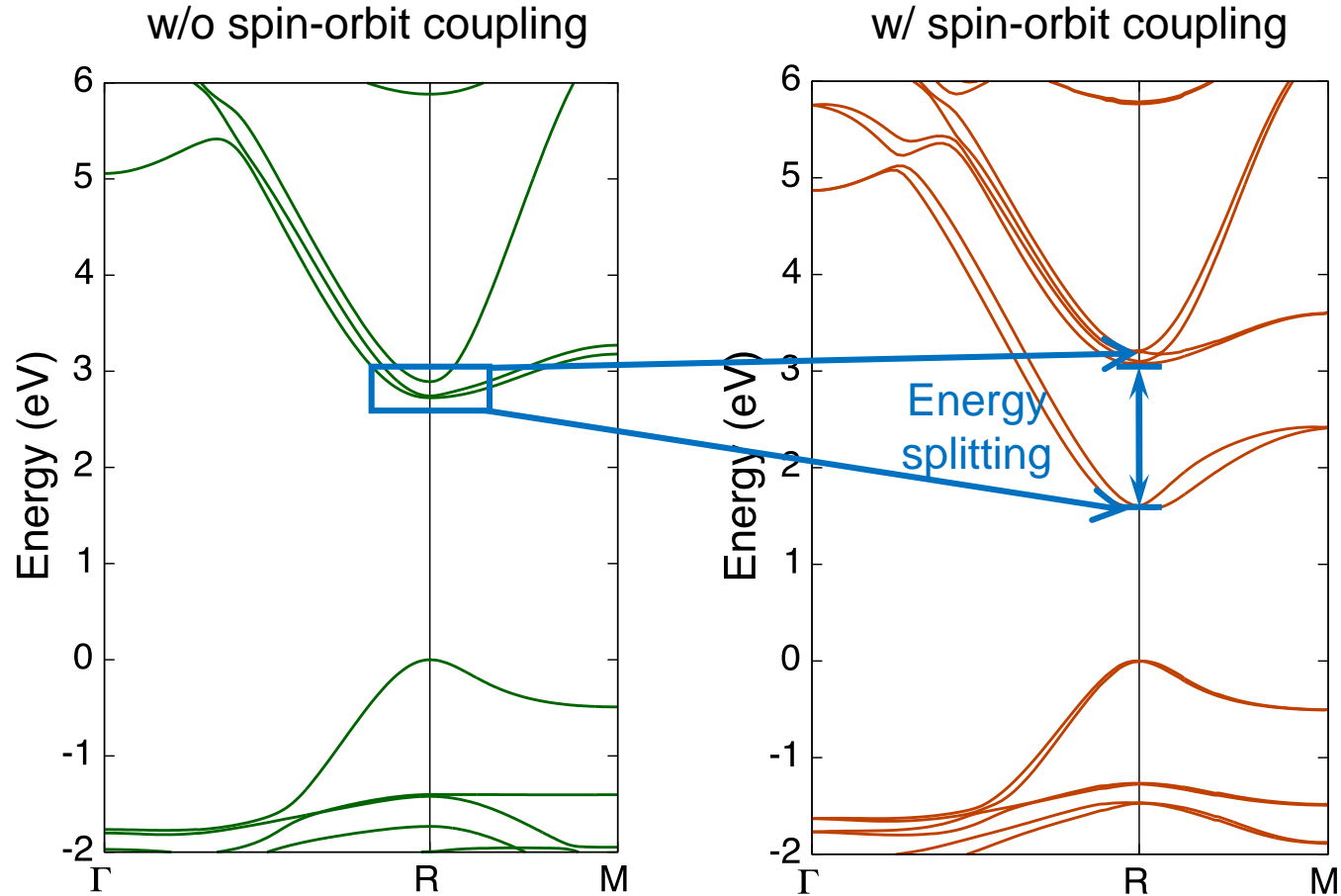
- MAPbI₃ has a much greater Auger coefficient $\sim 100\times$

Origin of strong Auger



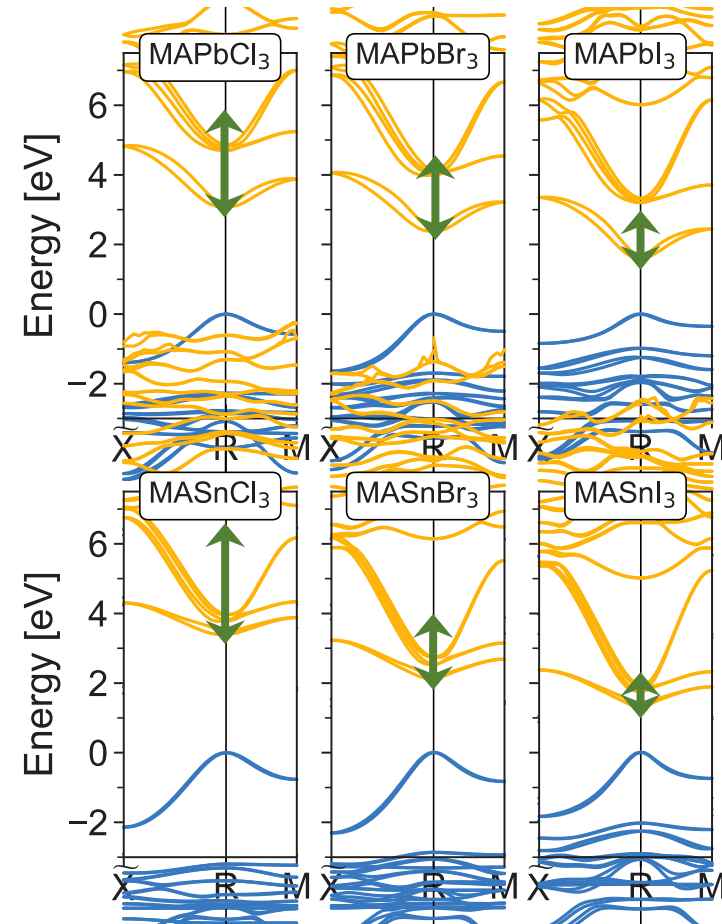
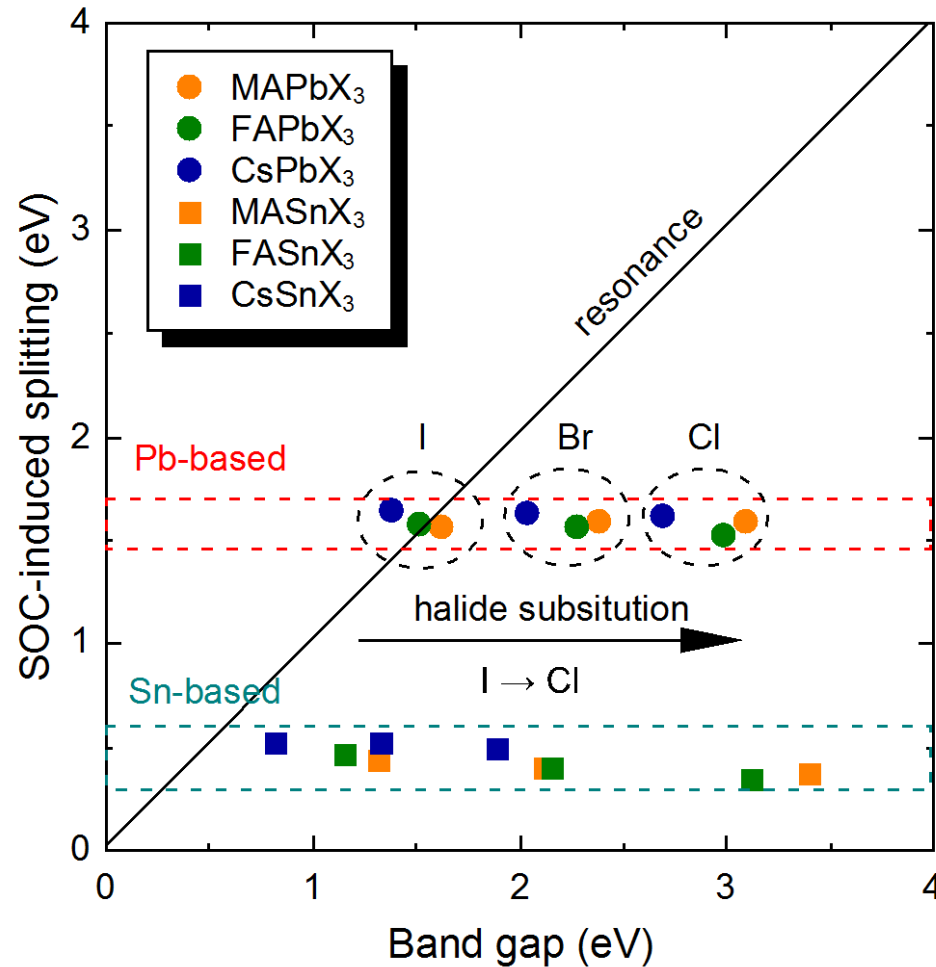
J.-X. Shen, X. Zhang, S. Das, E. Kioupakis, and C. G. Van de Walle,
Adv. Energy Mater. **8**, 1801027 (2018).

Spin-orbit energy splitting



Energy splitting in conduction bands causes resonance for eeh Auger

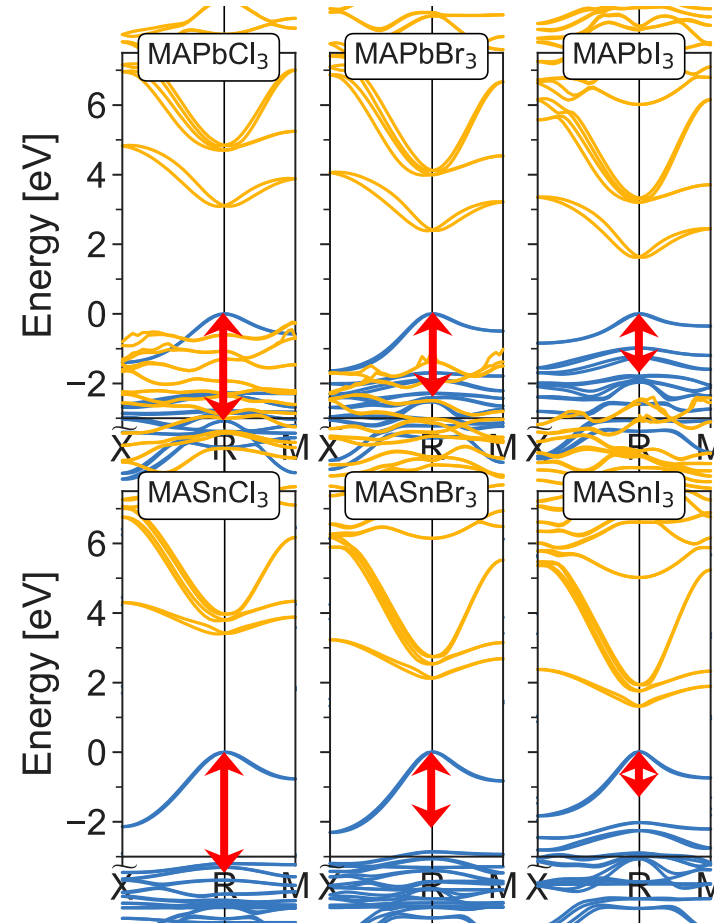
Band-structure engineering



X. Zhang, J.-X. Shen, and C. G. Van de Walle, *Adv. Energy Mater.* **9**, 1902830 (2019).

- X-site substitution can suppress eeh Auger

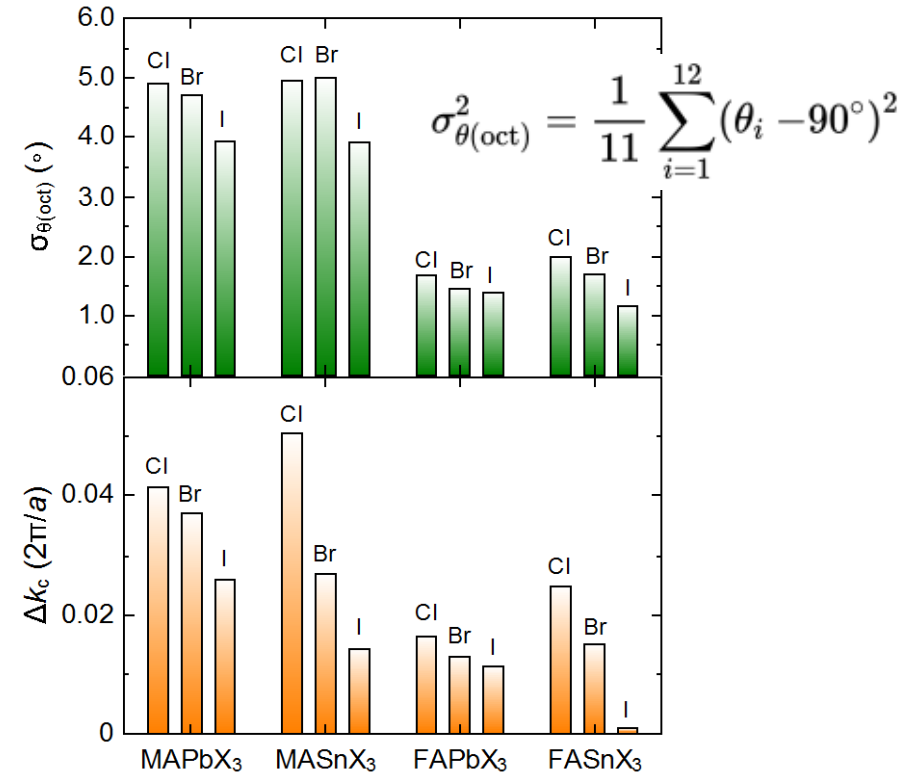
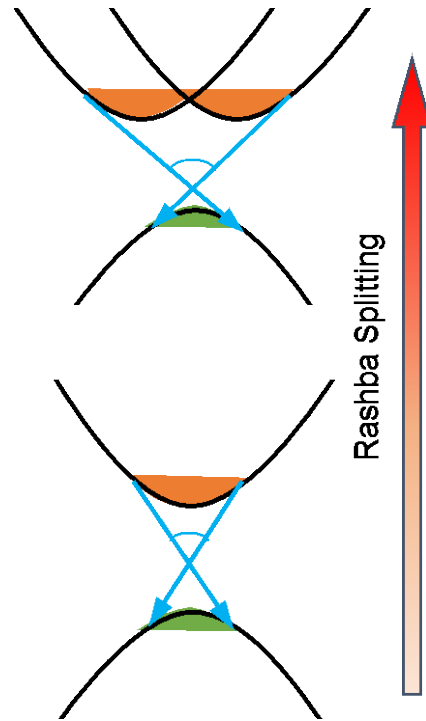
What about hhe Auger?



X. Zhang, J.-X. Shen, and C. G. Van de Walle, *Adv. Energy Mater.* **9**, 1902830 (2019).

- B-site substitution allows suppressing hhe Auger

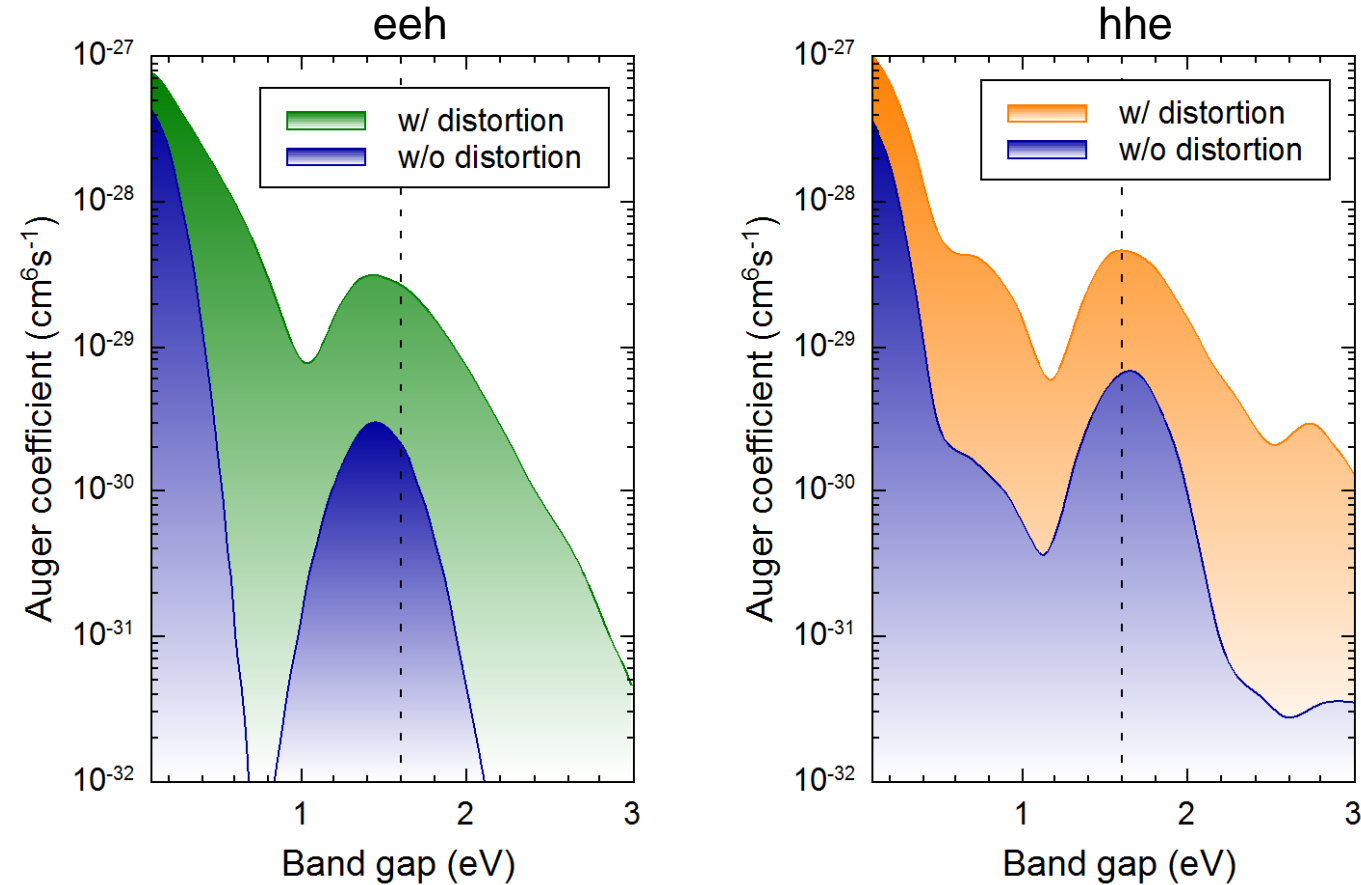
Suppressing lattice distortions



X. Zhang, J.-X. Shen, and C. G. Van de Walle, *Adv. Energy Mater.* **9**, 1902830 (2019).

- Tunable lattice distortion and Rashba splitting by chemical substitution

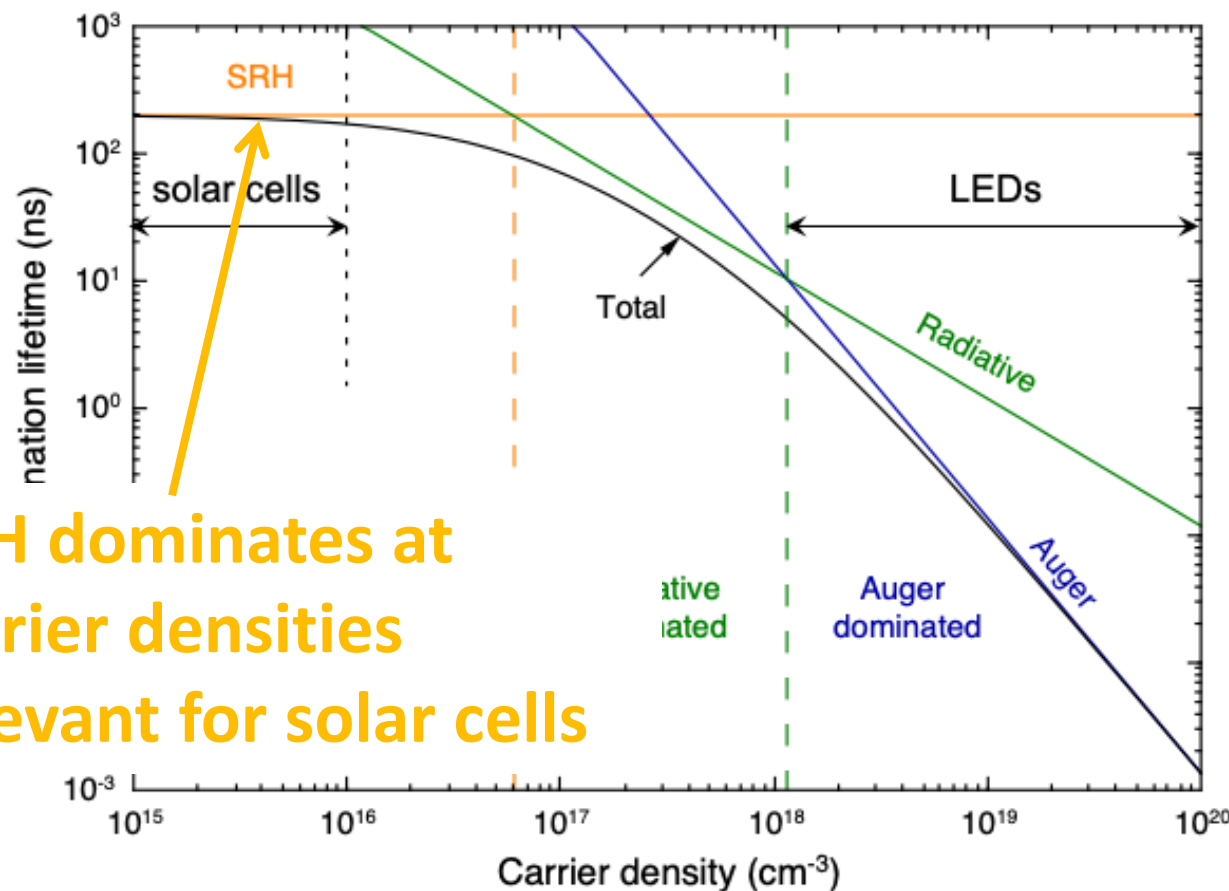
Suppressing lattice distortions



- Suppressing lattice distortions and thus the Rashba splitting reduces Auger by one order of magnitude

Recombination in halide perovskites

First-principles studies of recombination rates



SRH dominates at carrier densities relevant for solar cells

X. Zhang *et al.*, Adv. Energy Mater. **10**, 1902830 (2020).

X. Zhang *et al.*, J. Phys. Chem. Lett. **9**, 2903 (2018).

X. Zhang *et al.*, ACS Energy Lett. **3**, 2329 (2018).

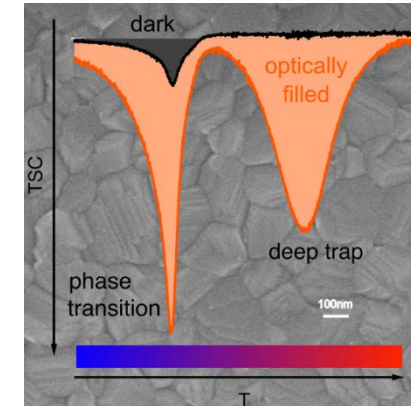
J.-X. Shen *et al.*, Adv. Energy Mater. **8**, 1801027 (2018).

- Defect-assisted (“SRH”) recombination limits efficiency

“Defect tolerance”

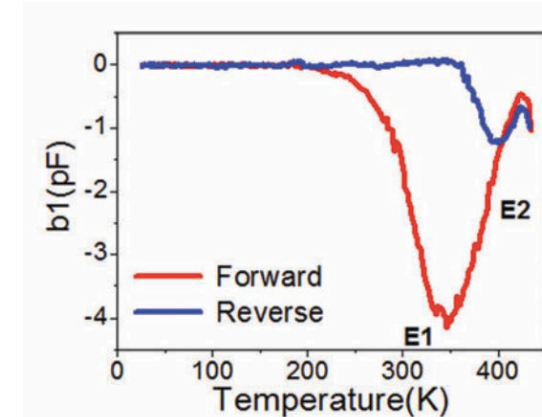
- “Defect tolerance”: defects are present, but do not cause strong nonradiative carrier recombination
- Concept emerged from early, less accurate, first-principles calculations for defects in MAPbI_3
 - none of the relevant defects had levels deep in the band gap
- Commonly invoked to explain the high efficiency of perovskite solar cells
- However, deep-level defects with concentrations $\sim 10^{15} \text{ cm}^{-3}$ are observed experimentally!

- Thermally stimulated current (TSC)



A. Baumann *et al.*, *J. Phys. Chem. Lett.* **6**, 2350 (2015).

- Deep level transient spectroscopy (DLTS)



S. Heo *et al.*, *Energy Environ. Sci.* **10**, 1128 (2017).

Point defects in halide perovskites

First-principles calculations of formation energies and defect levels

C. Freysoldt *et al.*, Rev. Mod. Phys. **86**, 253 (2014).

Example: iodine interstitial (I_i)

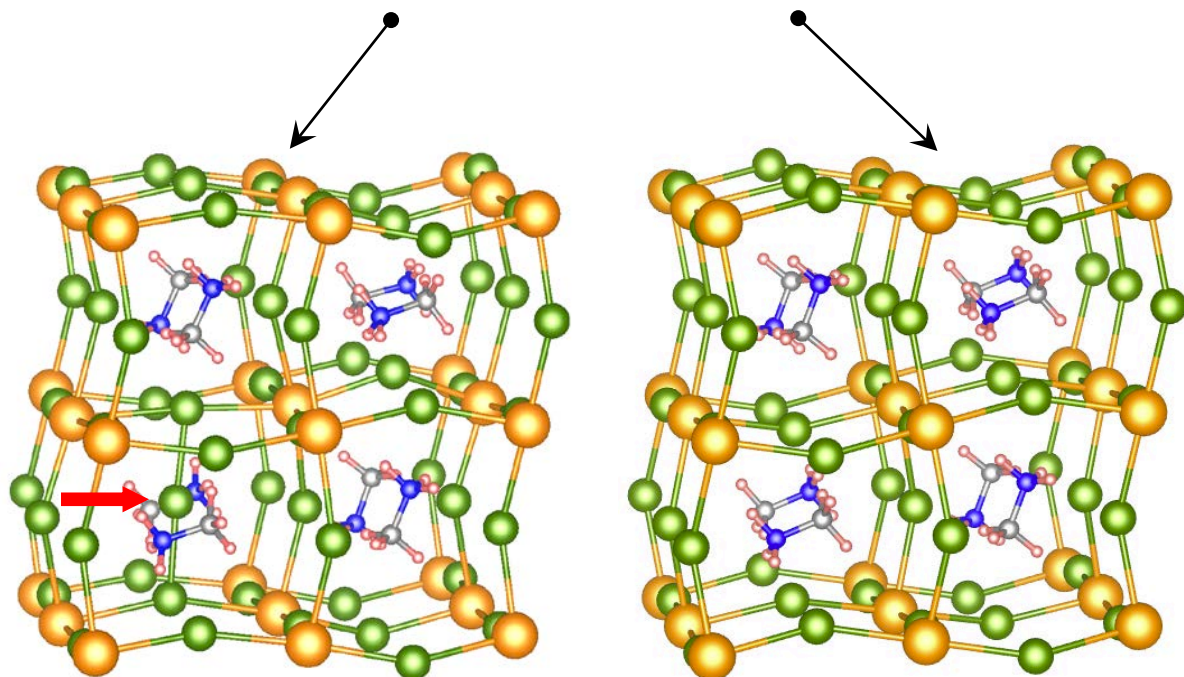
$$E^f [I_i^q] = E_{\text{tot}} [I_i^q] - E_{\text{tot}} [\text{bulk}] - \mu_I + qE_F$$

● (chemical potential of I)

Fermi level
(chemical potential of electrons)

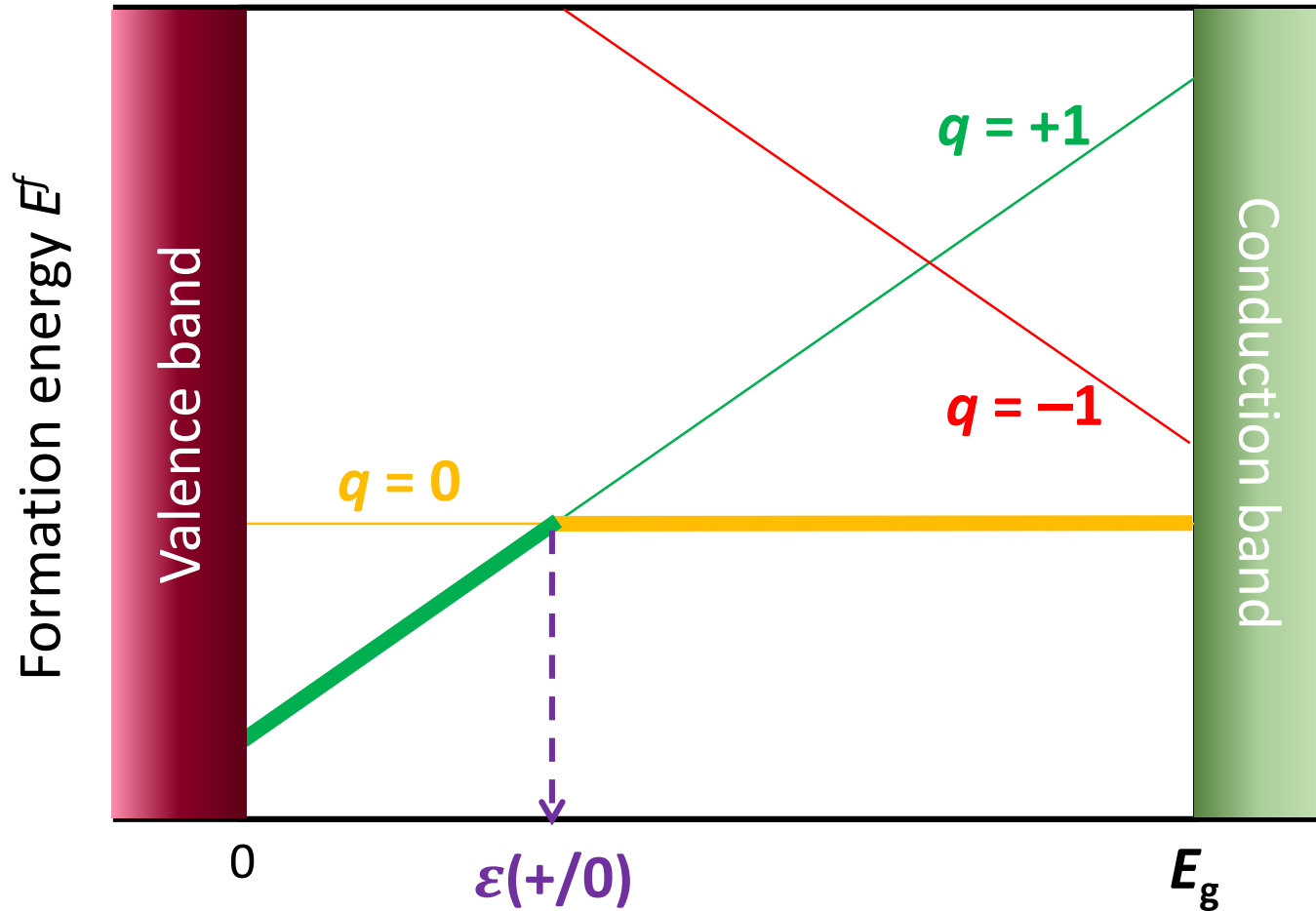
- Density functional theory
- Hybrid functional
 - J. Heyd, G. E. Scuseria, and M. Ernzerhof, J. Chem. Phys. **118**, 8207 (2003).
- Spin-orbit coupling
- Supercells, atomic relaxation

Defect concentration: $N_{\text{def}} = N_{\text{sites}} e^{-E^f/k_B T}$

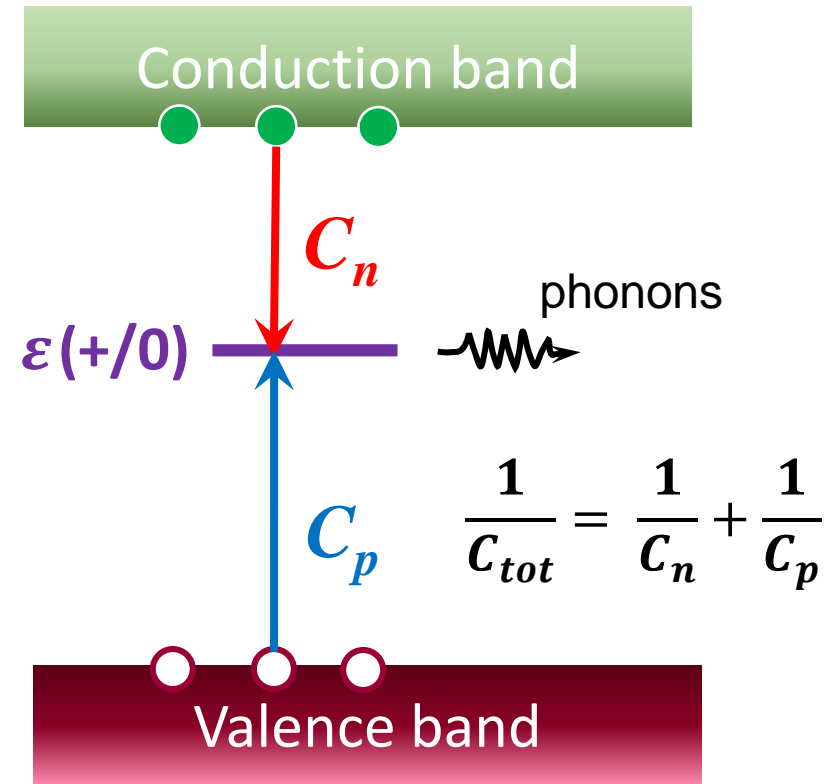


$q=+1$

Defect-assisted recombination in halide perovskites

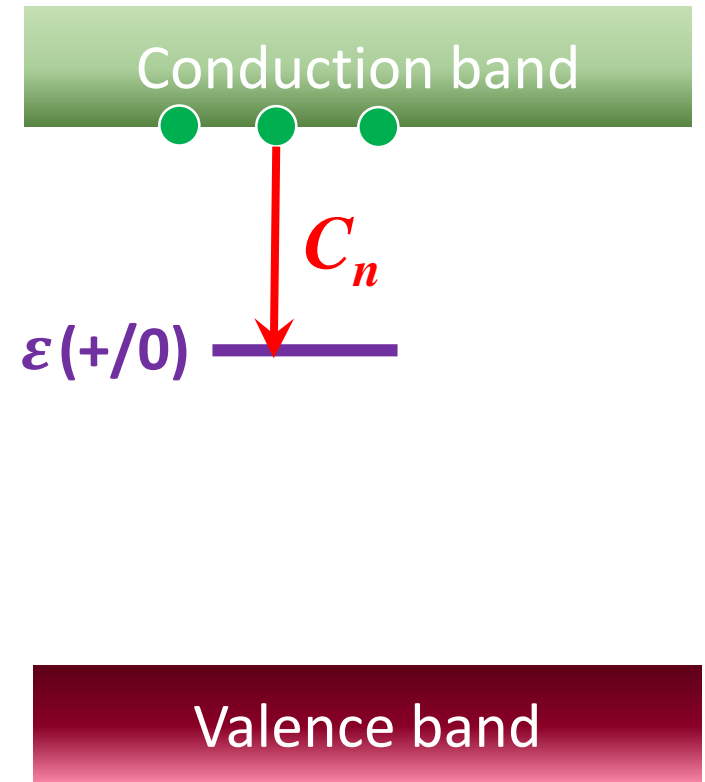
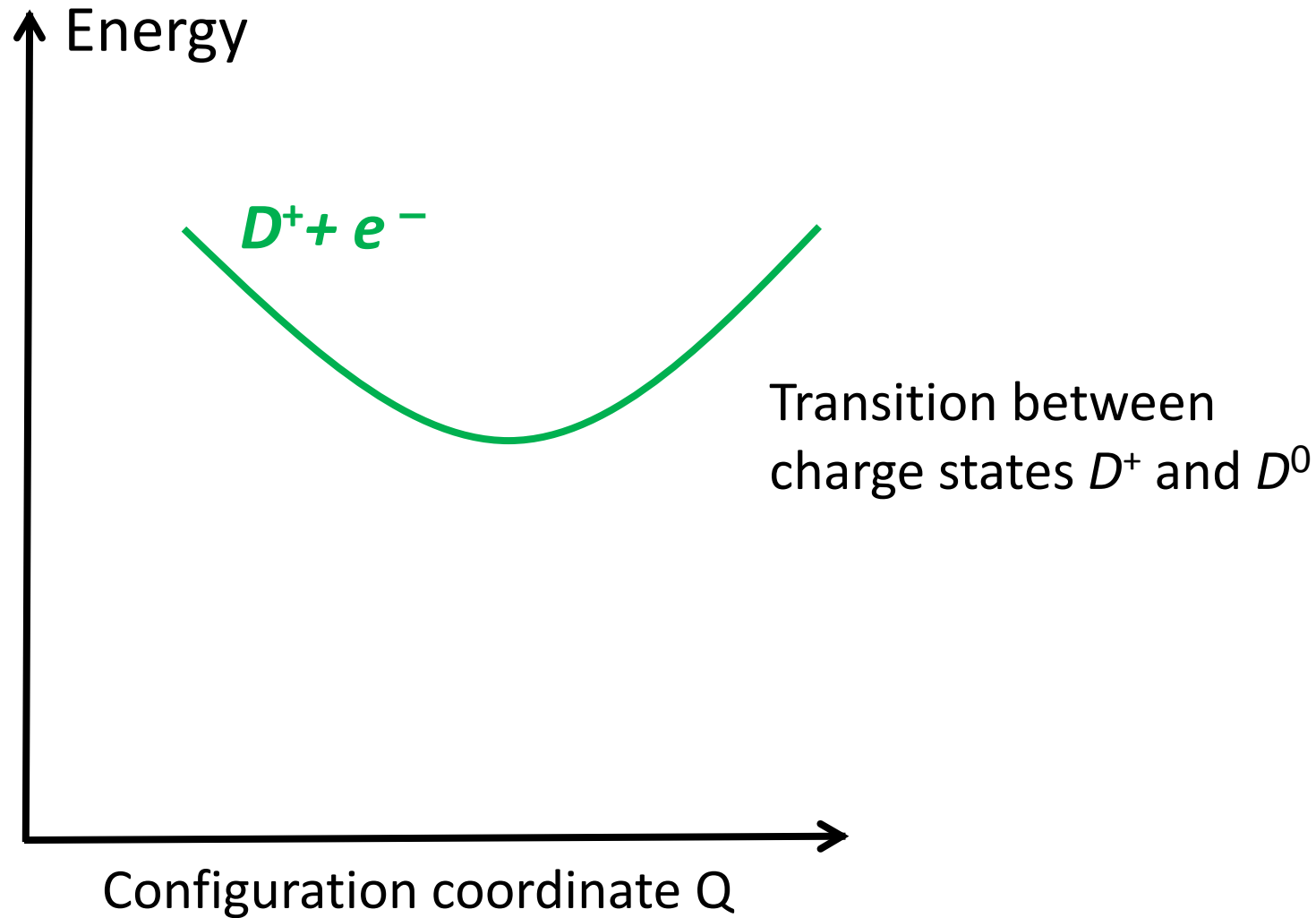


$\epsilon(+/0)$: charge-state transition level relevant for SRH recombination

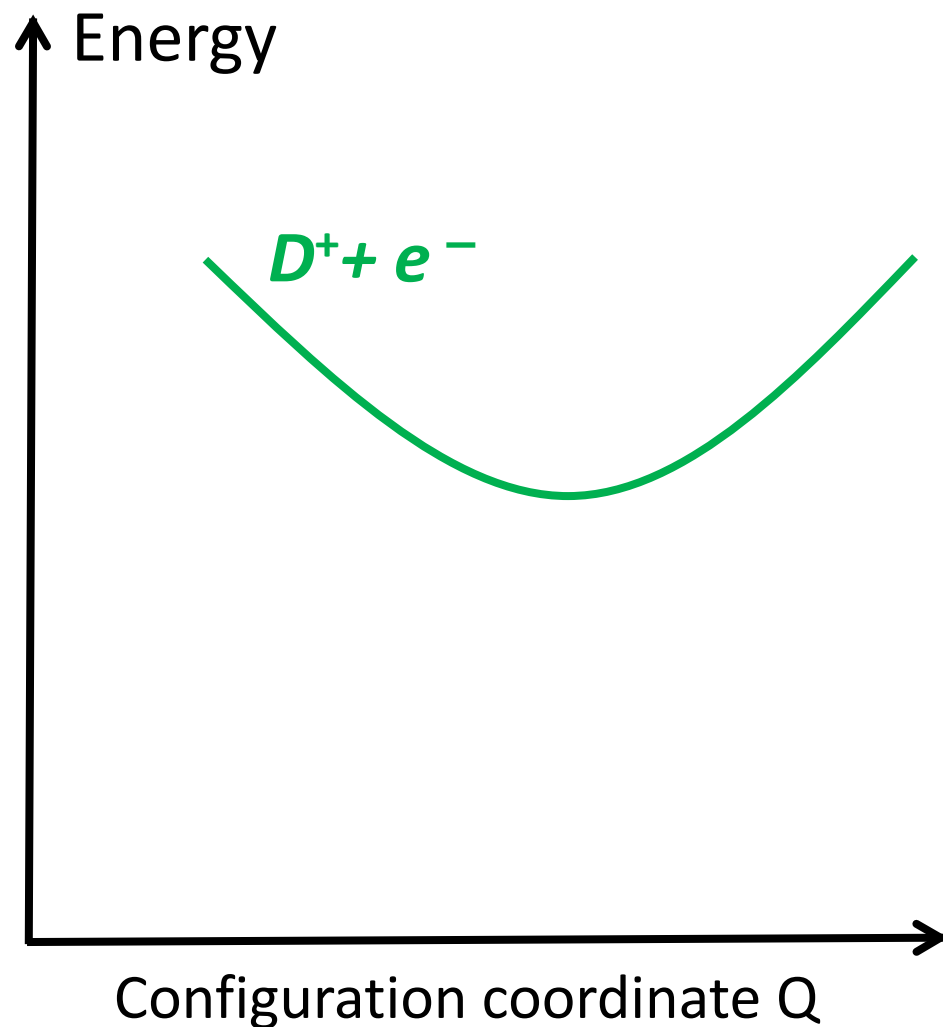


- Usually, $C \propto e^{-\Delta E}$
- Mid-gap defects have best balance between C_n and C_p

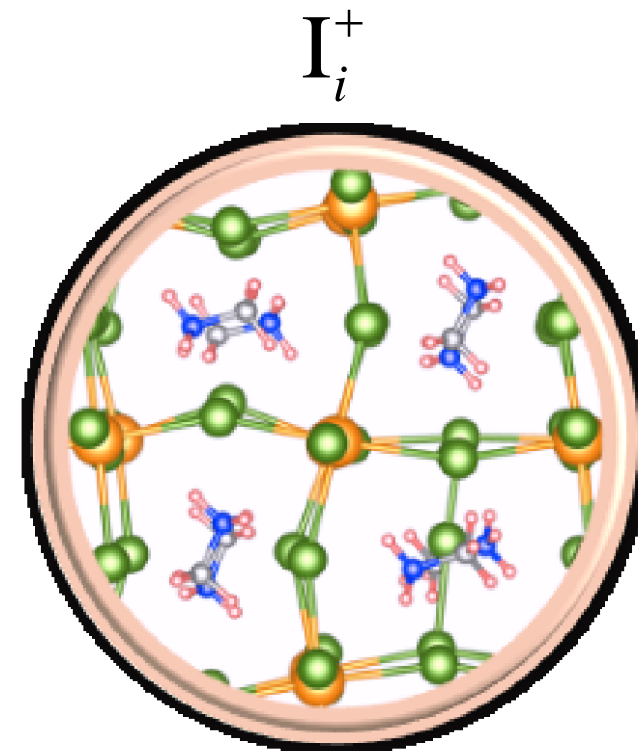
Configuration coordinate diagram



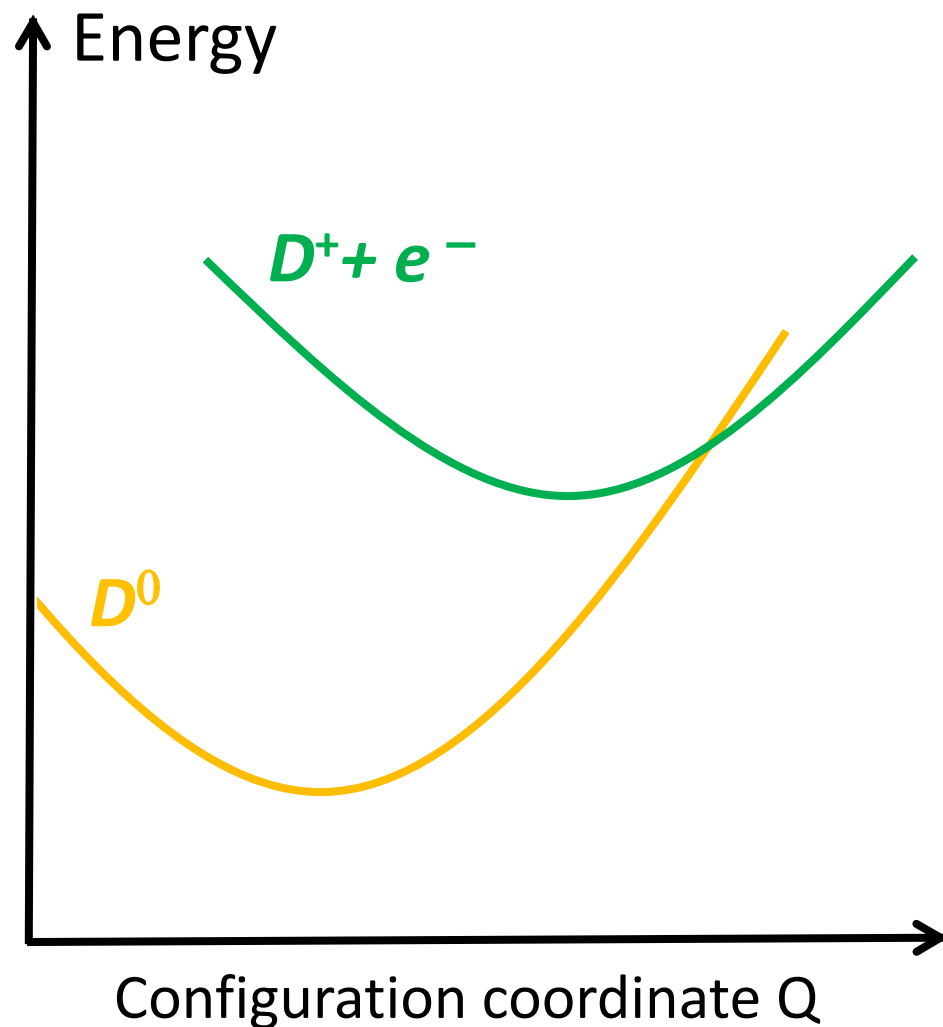
Configuration coordinate diagram



Energy depends on atomic configuration

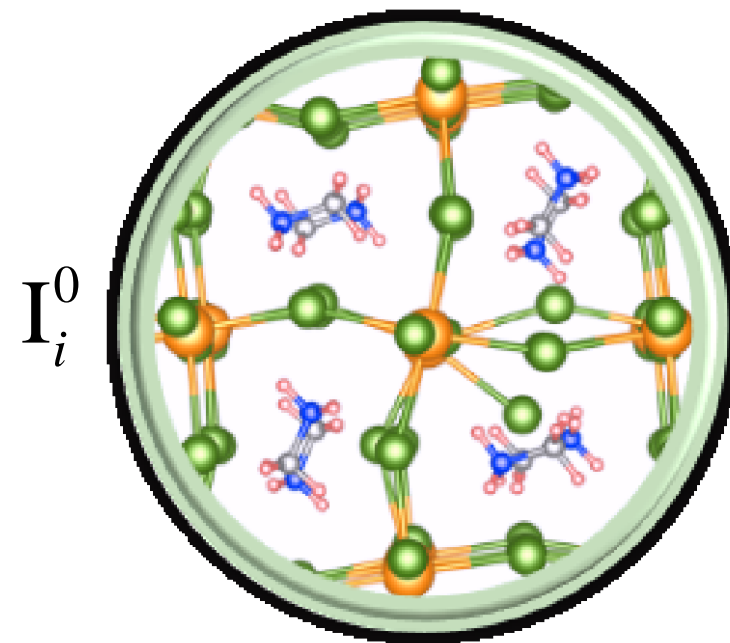


Configuration coordinate diagram

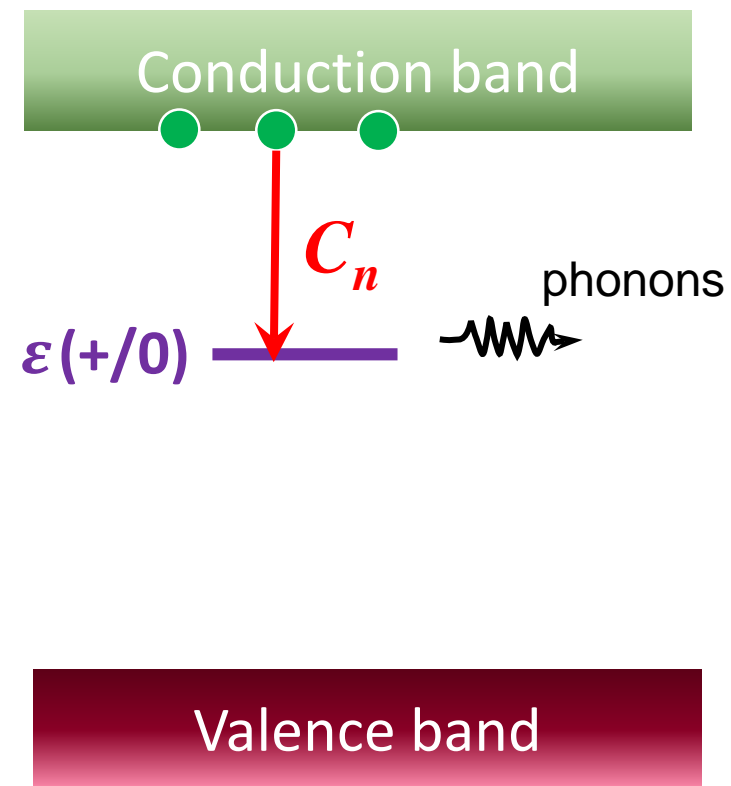
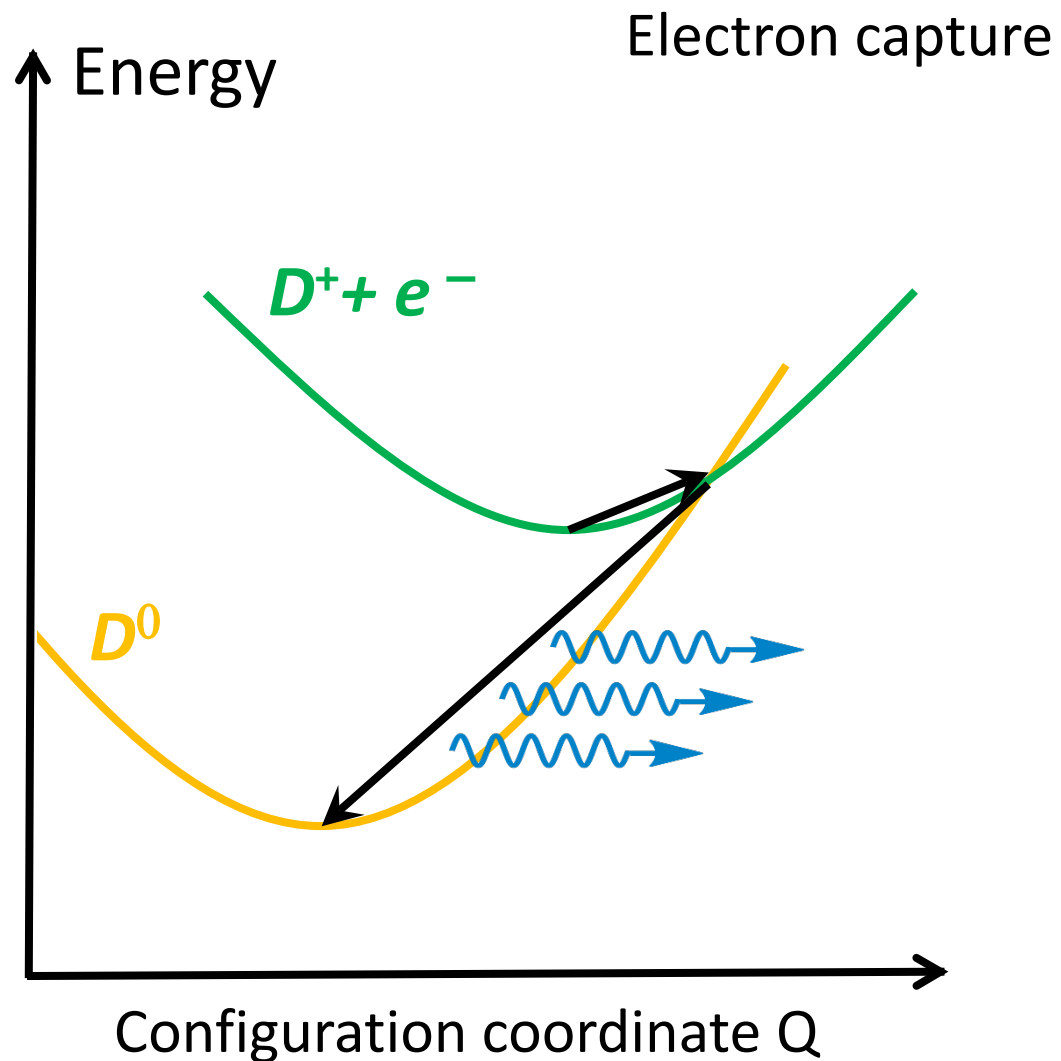


Energy depends on atomic configuration

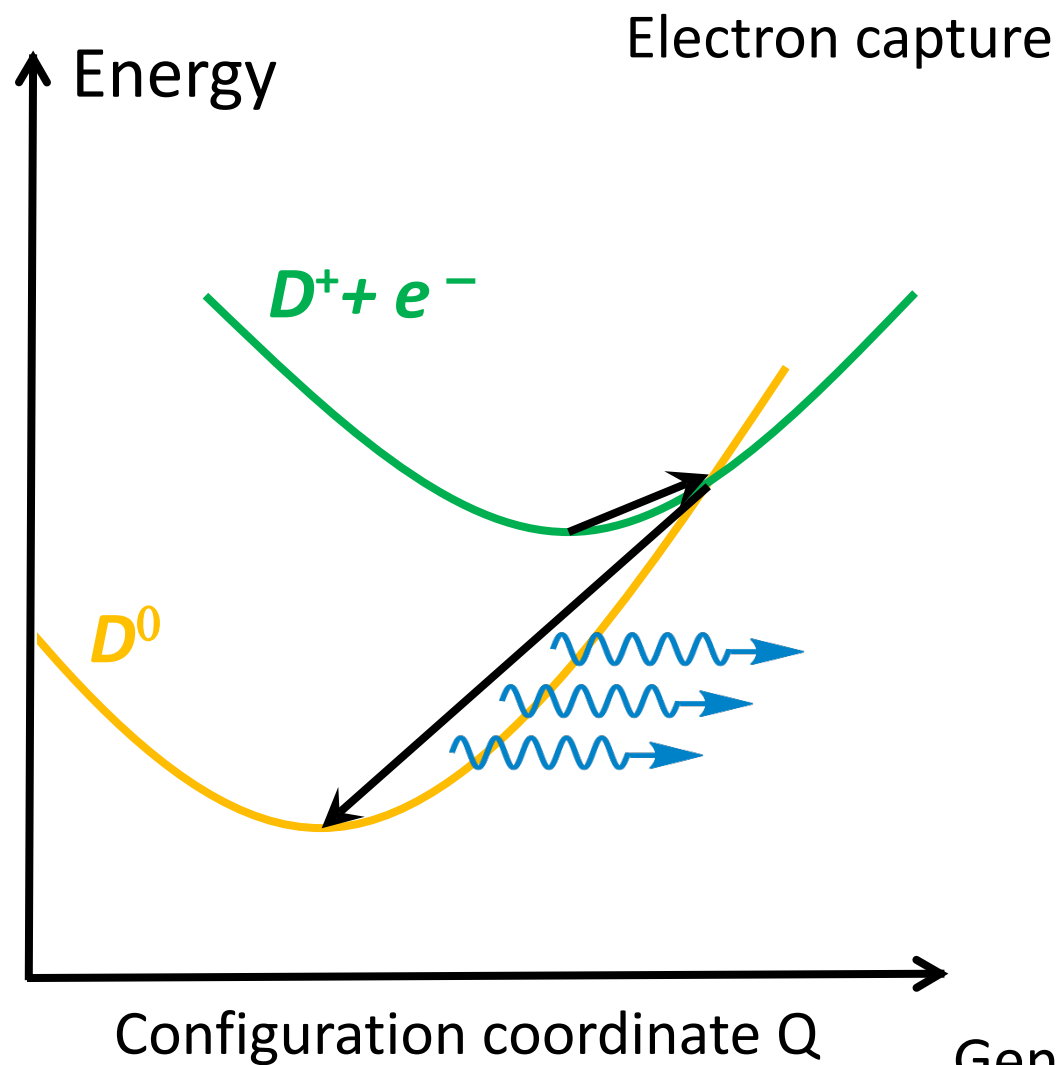
Charge state D^0 has different atomic configuration from D^+



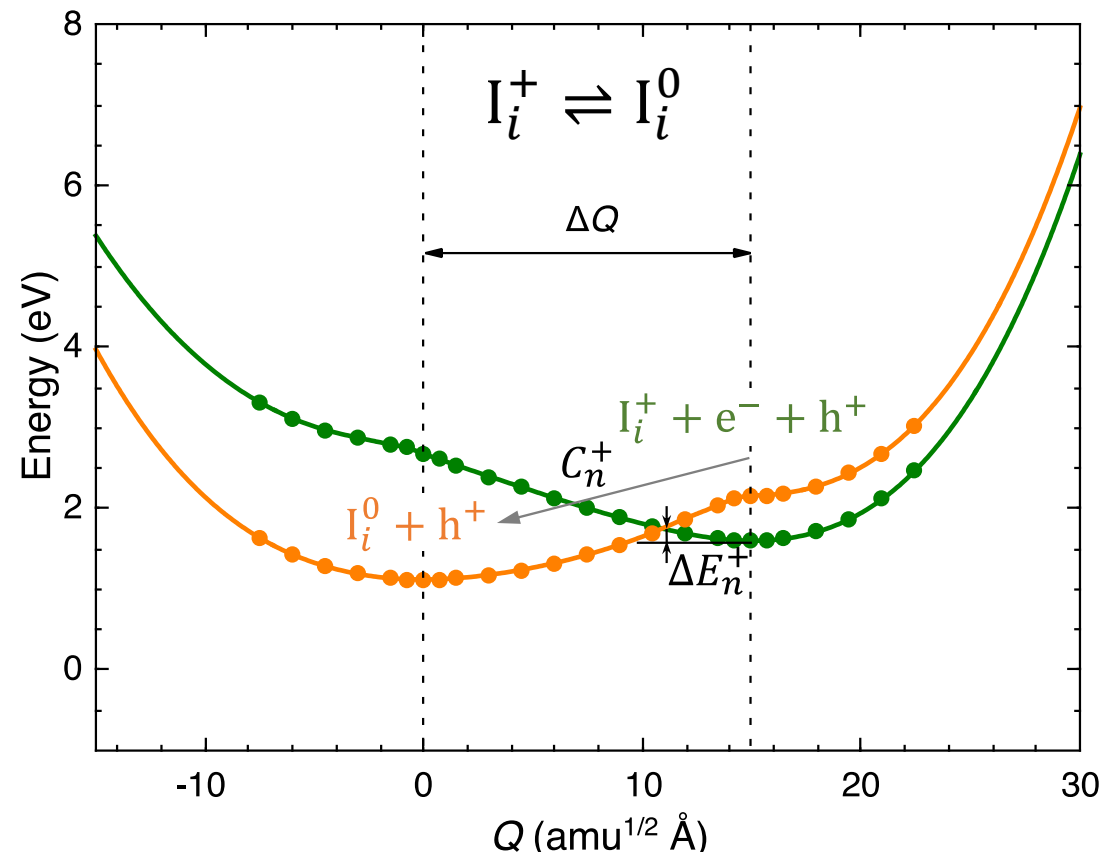
Configuration coordinate diagram



Configuration coordinate diagram



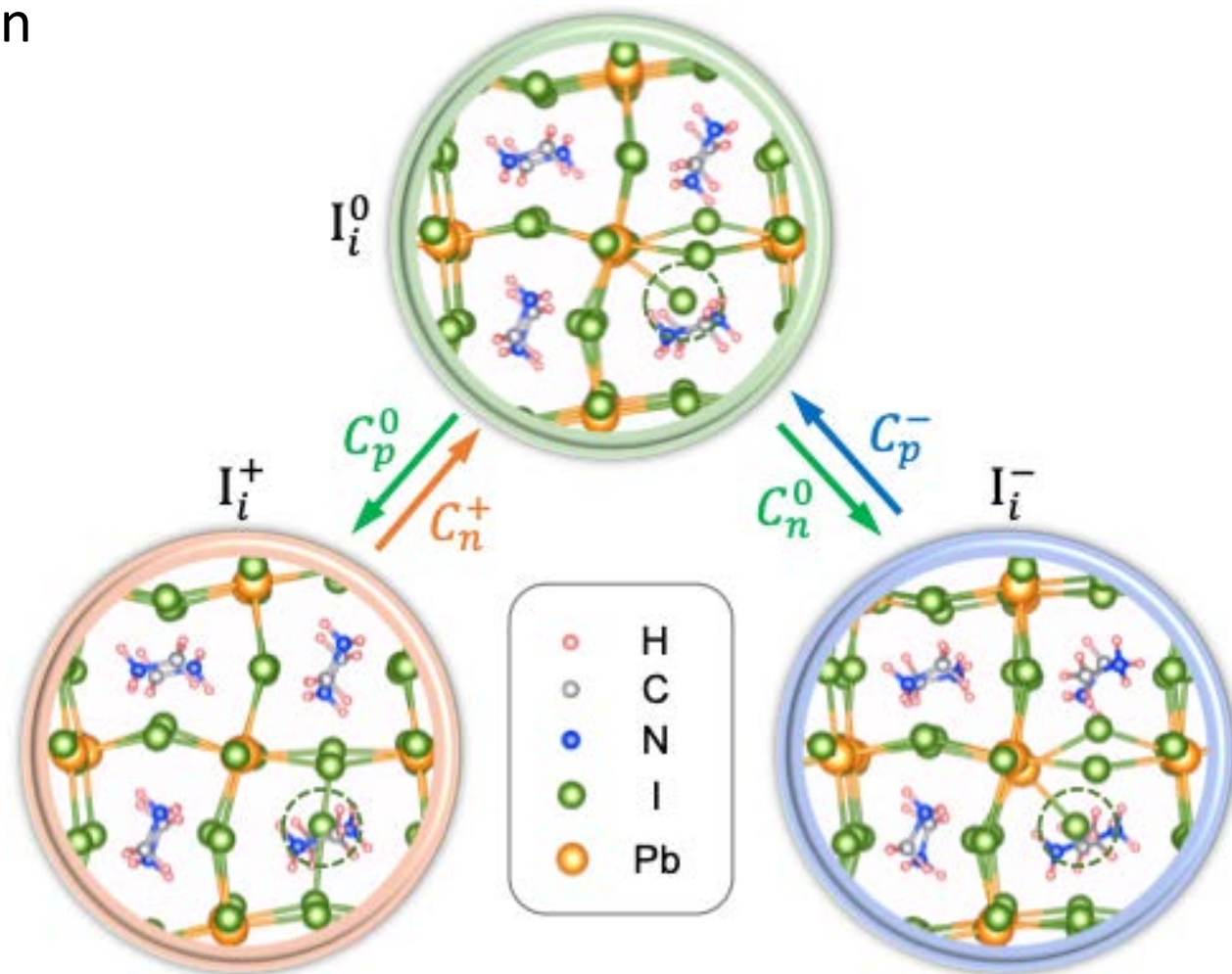
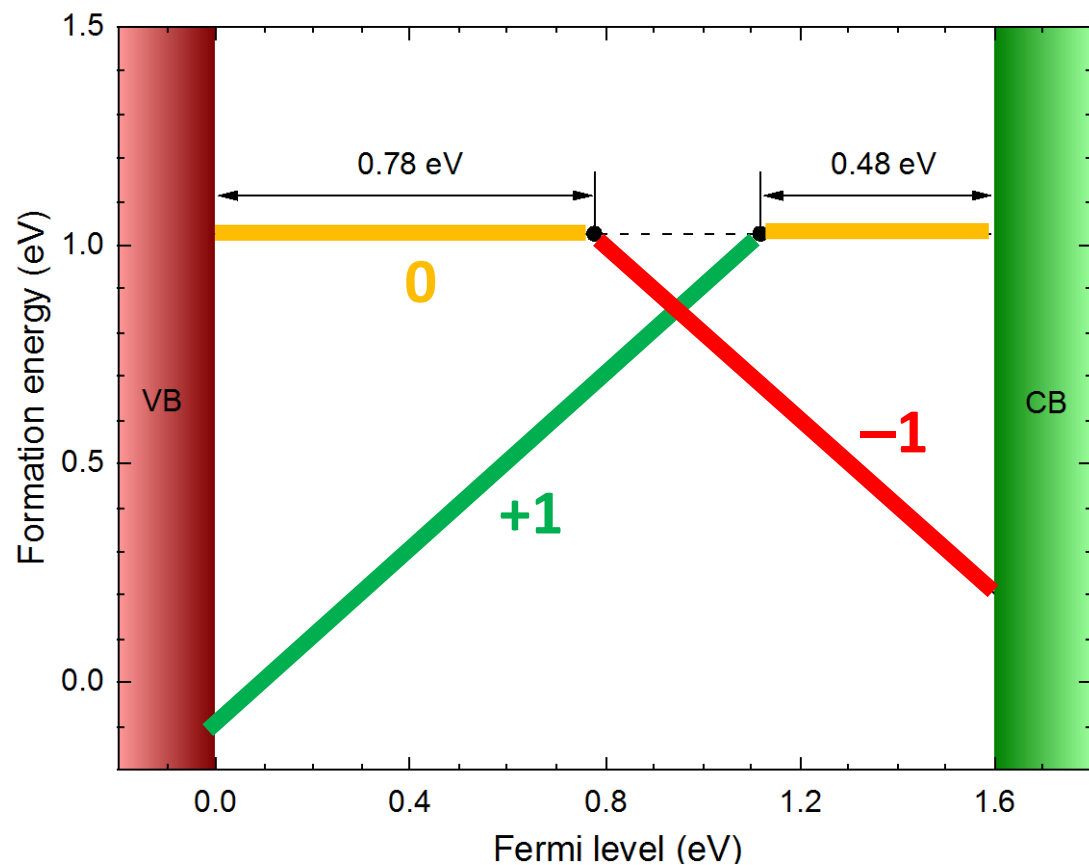
Configuration coordinate diagram from first principles



Generalized coordinate: $Q = \sqrt{\sum_{\alpha} m_{\alpha} (\mathbf{R}_{\alpha} - \mathbf{R}_{f;\alpha})^2}$

Prominent defect: iodine interstitial

- Low formation energy \rightarrow high concentration
- Four capture processes

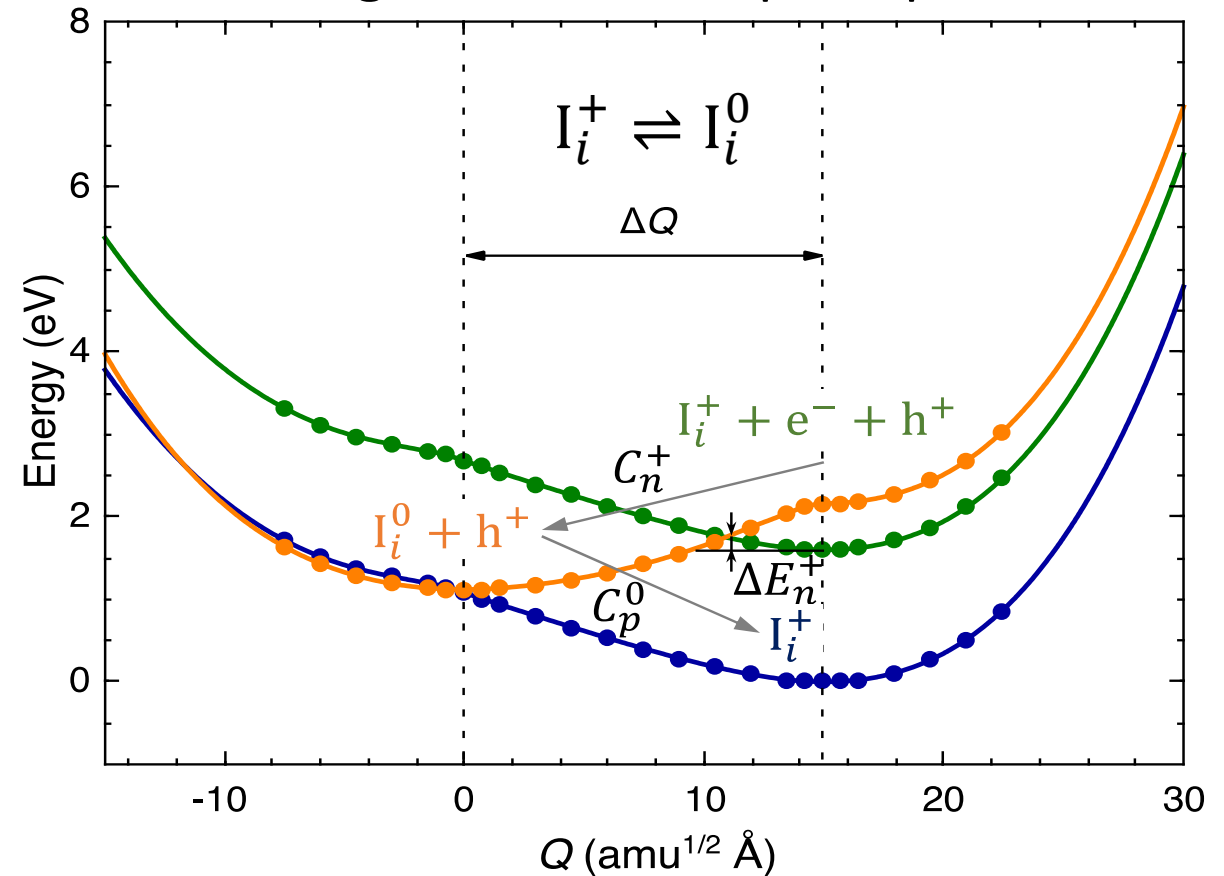


Configuration coordinate diagram: $I_i^+ \rightleftharpoons I_i^0$

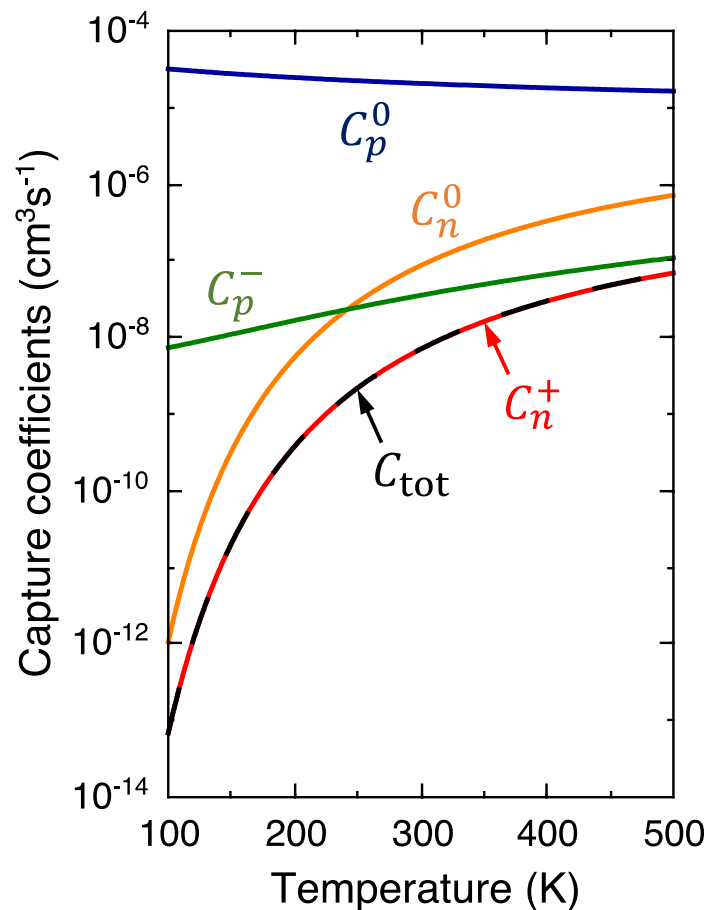
- Generally small capture barriers
→ high capture coefficients
- Capture coefficients do not decrease as expected with energy difference from band edge (i.e., $C \propto e^{-\Delta E}$)
- Two reasons:
 - Anharmonicity
 - “Marcus inverted region”

X. Zhang *et al.*, Phys. Rev. B **101**, 140101 (2020).

Configuration coordinate diagram from first principles

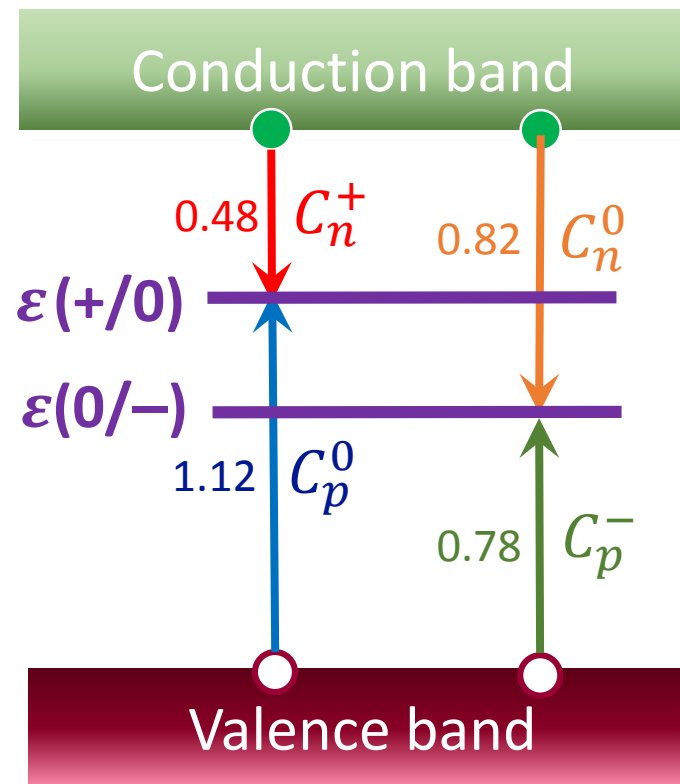


Capture coefficients



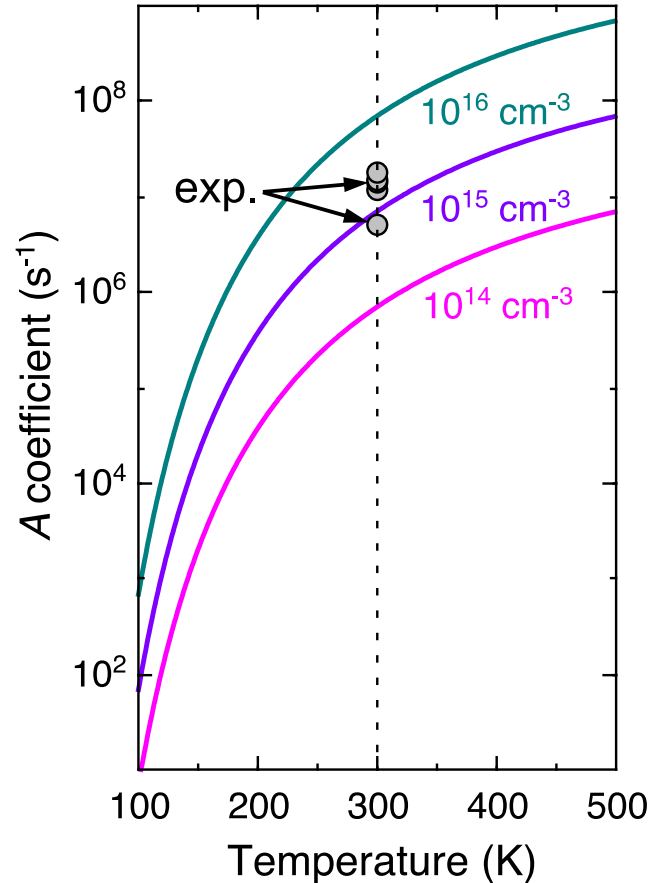
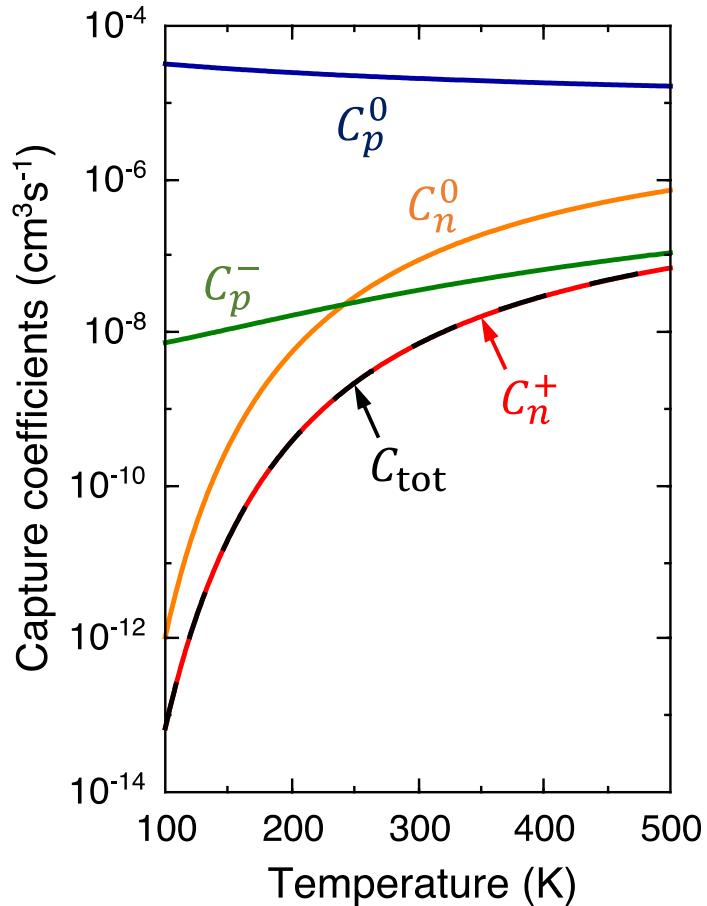
- $I_i^+ \rightleftharpoons I_i^0$ and $I_i^0 \rightleftharpoons I_i^-$ charge-state transitions
- Four capture processes
- Total capture coefficient:

$$C_{\text{tot}} = \frac{C_n^0 + C_p^0}{1 + \frac{C_n^0}{C_p^-} + \frac{C_p^0}{C_n^+}}$$



X. Zhang *et al.*, Phys. Rev. B **101**, 140101 (2020).

Capture coefficients



- Total capture coefficient:

$$C_{\text{tot}} = \frac{C_n^0 + C_p^0}{1 + \frac{C_n^0}{C_p^-} + \frac{C_p^0}{C_n^+}}$$

- Nonradiative recombination rate:

$$R = An ; A = N_{\text{def}} C_{\text{tot}}$$

- The iodine interstitial is an efficient nonradiative recombination center
- Likely responsible for the observed rates in experiments

- $N_{\text{def}} \sim 10^{15} \text{ cm}^{-3}$

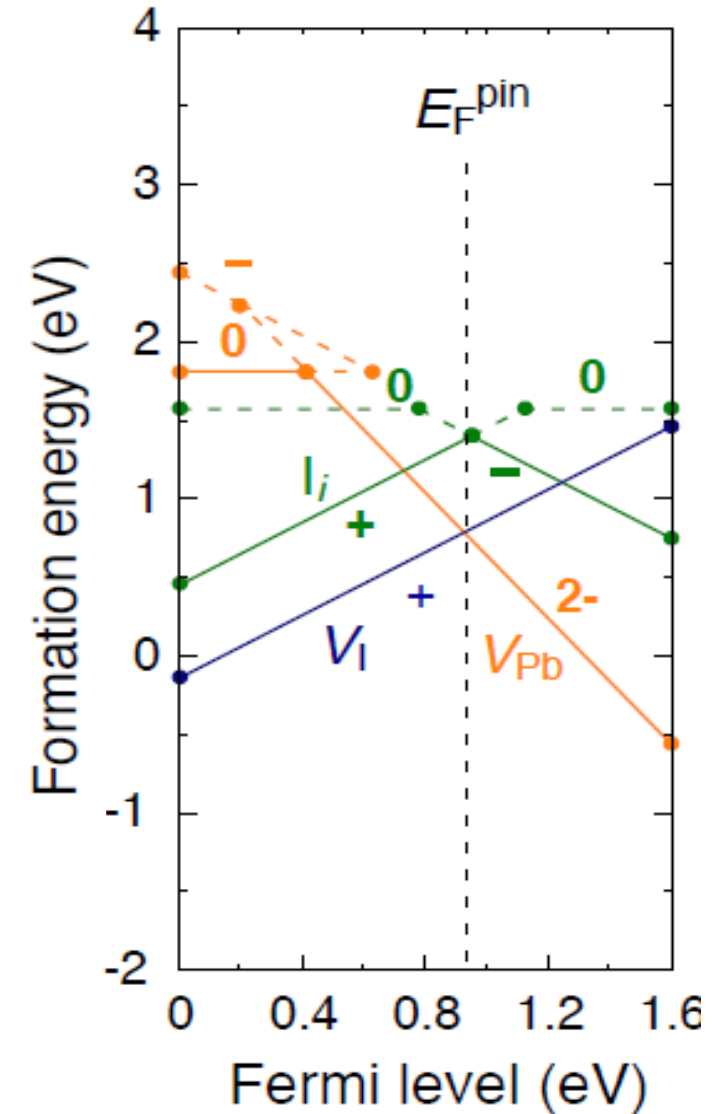
- A. Baumann *et al.*, J. Phys. Chem. Lett. **6**, 2350 (2015); S. Heo *et al.*, Energy Environ. Sci. **10**, 1128 (2017).

$$\Rightarrow A \approx 10^7 \text{ s}^{-1}$$

X. Zhang *et al.*, Phys. Rev. B **101**, 140101 (2020).

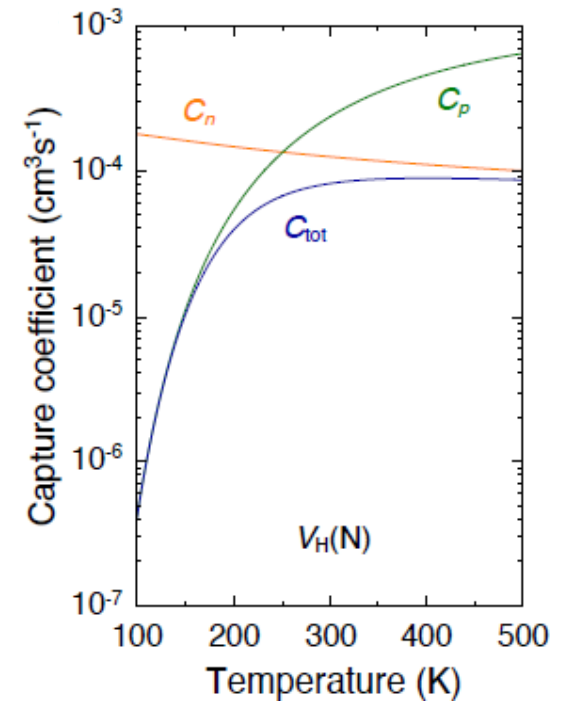
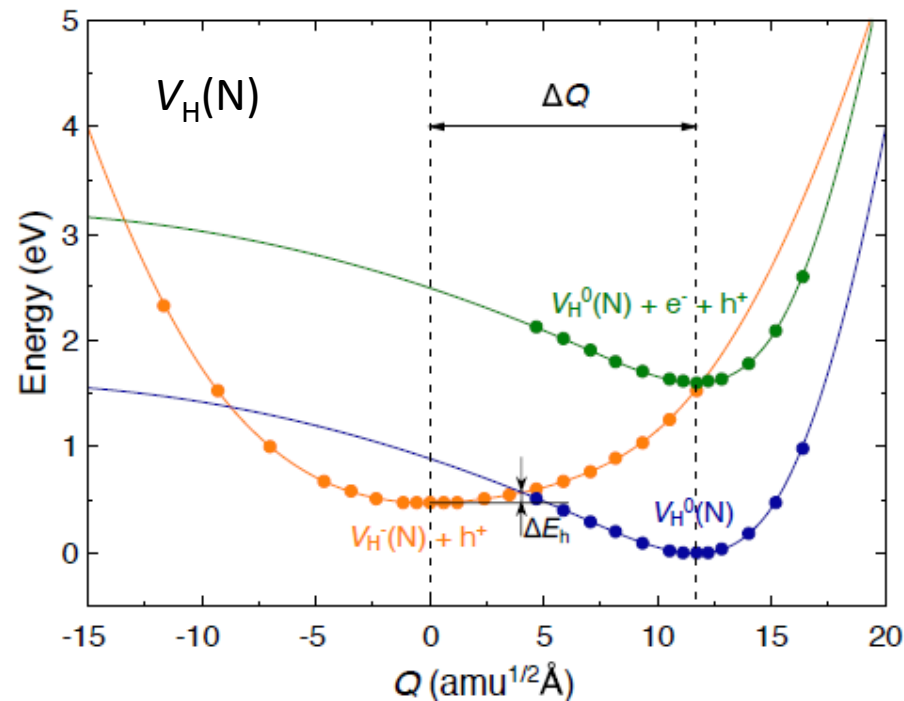
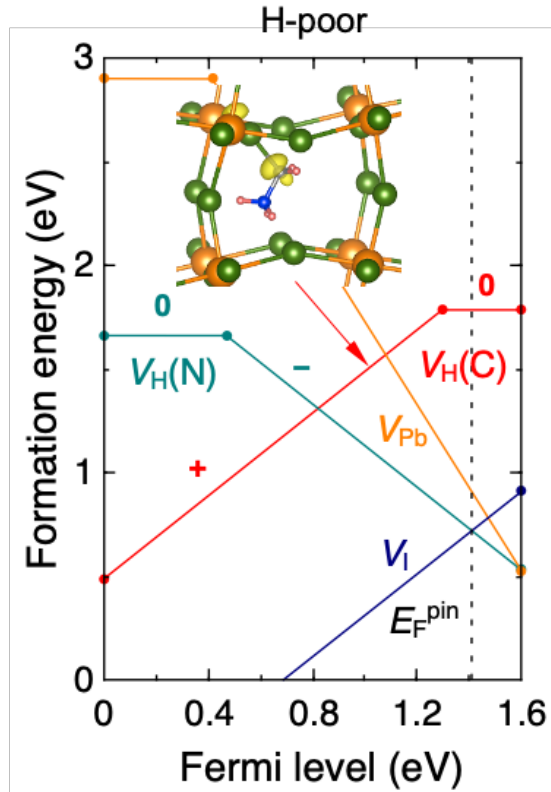
Other point defects

- Also examined other native point defects
 - **Pb interstitial (Pb_i) and antisites are high in energy**
⇒ unlikely to be present
 - **Iodine vacancy (V_I):** no charge-state transition levels in the band gap ⇒ cannot act as a recombination center
 - **Lead vacancy (V_{Pb}):** explicit calculations of recombination rates show that V_{Pb} does not cause efficient nonradiative recombination
- ⇒ **Iodine interstitial** likely responsible for the observed nonradiative recombination: $A \approx 10^7 \text{ s}^{-1}$ for $N_{\text{def}} \sim 10^{15} \text{ cm}^{-3}$
- Iodine-rich synthesis conditions should be avoided
 - Extreme iodine-poor should be avoided as well
 - promote the formation of hydrogen vacancies



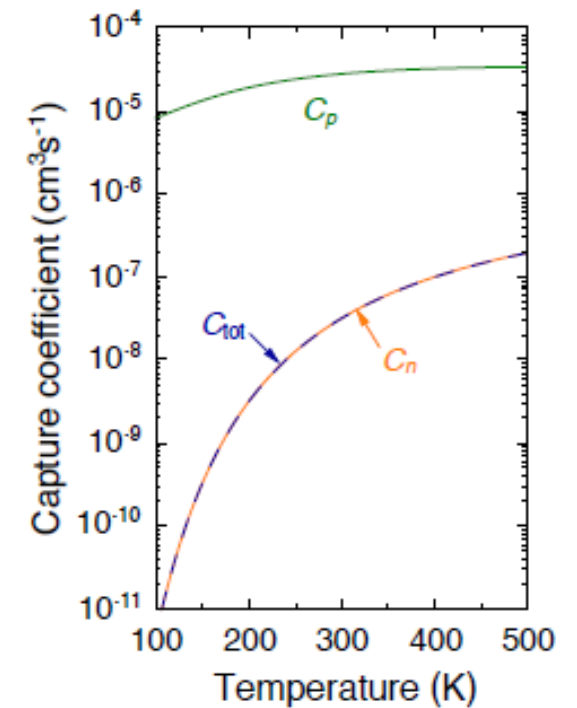
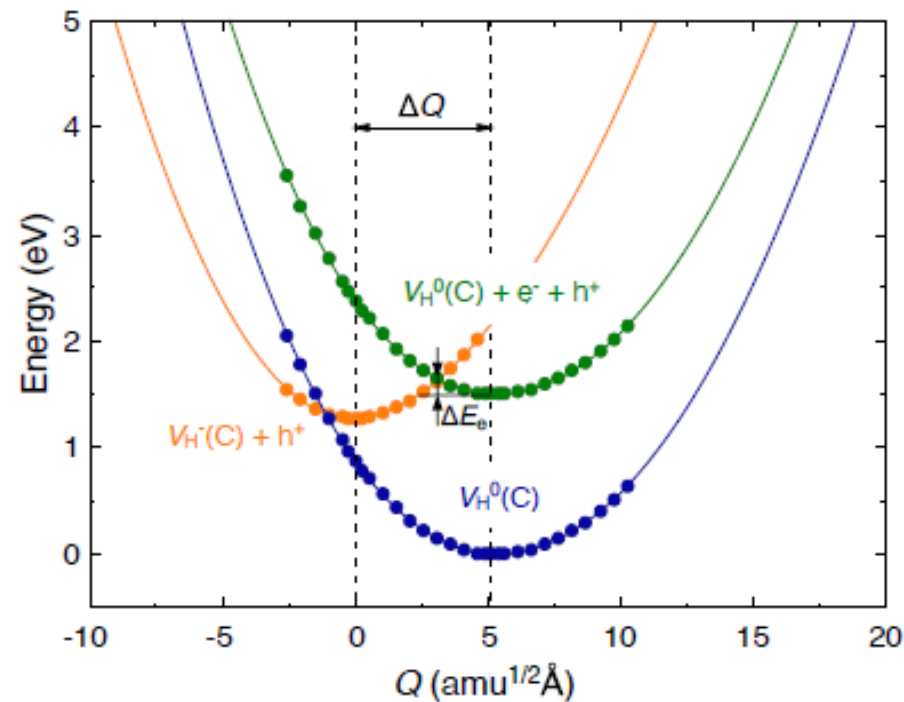
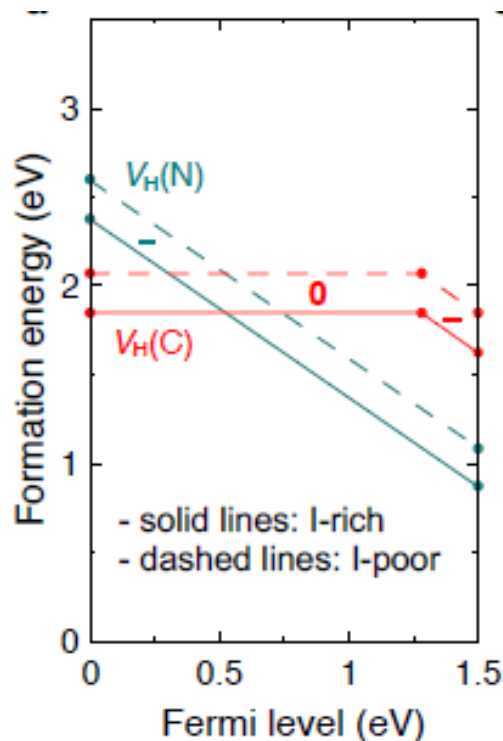
Hydrogen vacancies in MAPbI₃

- Hydrogen vacancies have been mostly overlooked when considering point defects
- MA: CH₃NH₃ → two types of H vacancies
 - V_H(C): removing H from a C atom
 - V_H(N): removing H from a N atom
- V_H(N) is an exceptionally strong recombination center (10⁻⁴ cm³s⁻¹)
- Present in high concentrations under I-poor and H-poor conditions



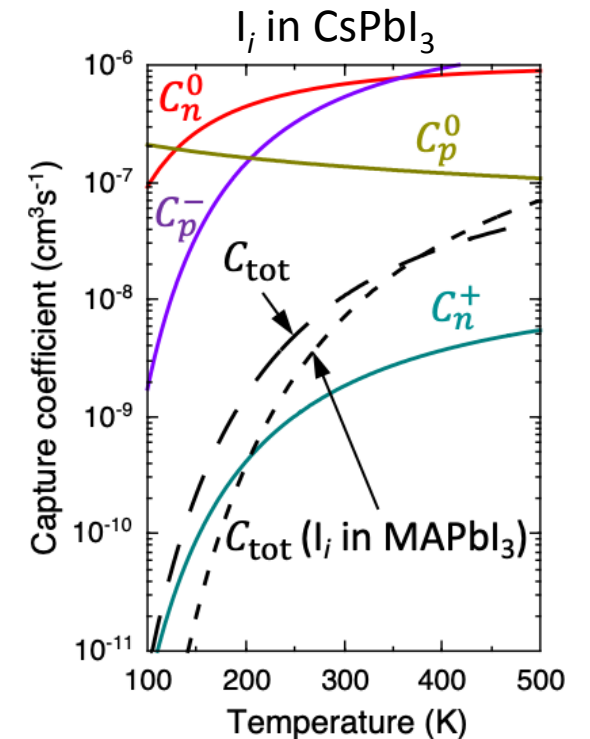
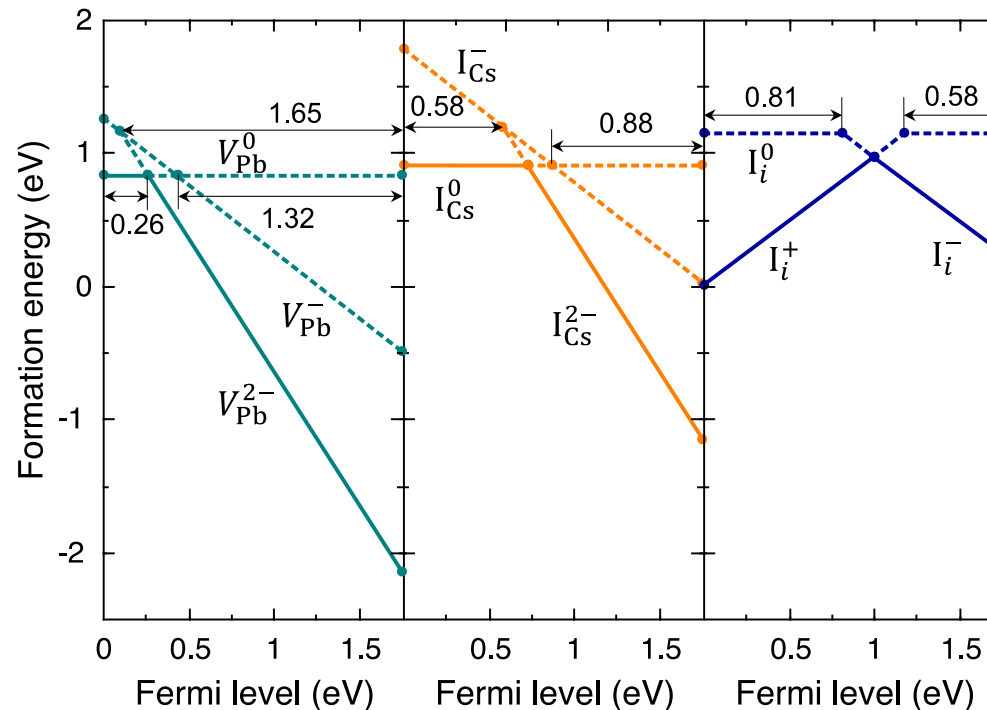
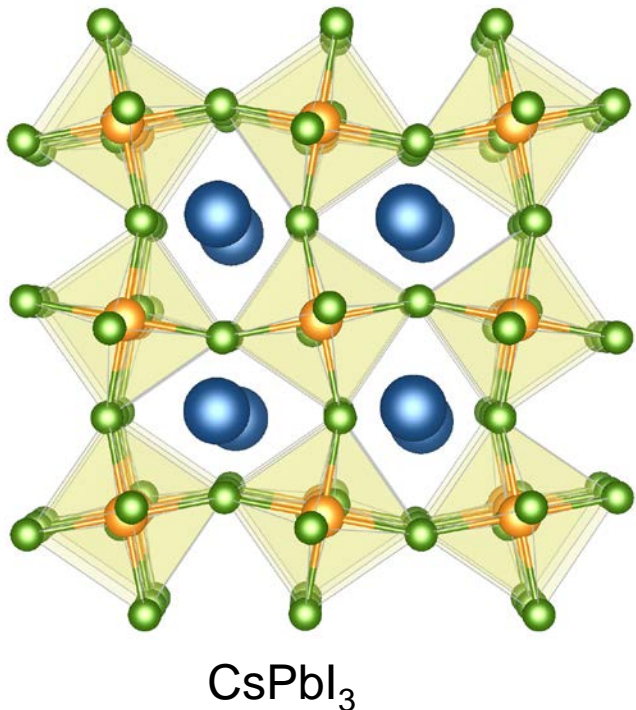
Qualitatively different behavior in FAPbI₃

- FA (formamidinium): CH(NH₂)₂
- $V_H(C)$ and $V_H(N)$ have much higher formation energies (lower concentrations) than in MAPbI₃
- $V_H(C)$ has substantially lower capture coefficient
- Rationalizes why FA is essential for realizing high efficiency



Getting rid of H vacancy problem by using CsPbI₃

- CsPbI₃: deep-level defects are present (V_{Pb} , I_{Cs} , and I_i)
- Explicit computation of recombination coefficients: I_i is the dominant recombination center
- Similar total capture coefficient as I_i in MAPbI₃, but no need to worry about H vacancies!
- Origin of current inferior performance of CsPbI₃: poor stability (small Cs⁺ → small tolerance factor)
- Ways to enhance stability: alloying, strain, and improved growth techniques



Putting values in perspective

- A coefficient in halide perovskites is **comparable to or higher than those in conventional semiconductors**
- Halide perovskites are often called **“defect tolerant”**
 - “Defects may be present, but do not harm efficiency”
 - Notion based on older calculations showing that defects do not introduce deep levels in the band gap
- Calling hybrid perovskites **“defect tolerant”** is misleading
 - We do not call GaAs or GaN “defect tolerant” —we worry a great deal about defects!
- Distinctive feature of hybrid perovskites: they can be grown with low defect densities using low-cost deposition techniques

Material	A coefficient (s ⁻¹)
MAPbI ₃	1.4 × 10 ⁷
MAPbI ₃	1.5 × 10 ⁷
MAPbI _{3-x} Cl _x	0.5 × 10 ⁷
MAPbI _{3-x} Cl _x	1.2 × 10 ⁷
FAPbI ₃	0.7 × 10 ⁷
FAPbBr ₃	2.1 × 10 ⁷
GaN	0.1 – 1.0 × 10 ⁷
GaAs	0.05 – 0.4 × 10 ⁷

M. B. Johnston *et al.*, *Acc. Chem. Res.* **49**, 146 (2016).
F. Olivier *et al.*, *Appl. Phys. Lett.* **111**, 022104 (2017).
E. Yablonovitch *et al.*, *Appl. Phys. Lett.* **50**, 1197 (1987).

Summary

- Rigorous first-principles calculations **elucidate mechanisms**
- **Radiative recombination**
 - Rashba spin splitting: spin texture, normal optical transitions
 - Rashba momentum splitting: limited impact, a factor of 2
J. Phys. Chem. Lett. **9**, 2903 (2018); ACS Energy Lett. **3**, 2329 (2018).
- **Auger recombination**
 - Resonance in band structure
 - Band-structure and lattice-distortion engineering allows reducing Auger
Adv. Energy Mater. **8**, 1801027 (2018); Adv. Energy Mater. **9**, 1902830 (2019).
- **Defect-assisted SRH recombination**
 - Halide perovskites often touted as “defect tolerant”; our work demonstrates that **defects do impact efficiency**.
 - Hydrogen-related defects act as strong nonradiative recombination centers
X. Zhang *et al.*, Phys. Rev. B **101**, 140101 (2020);
J. Phys. Chem. C **124**, 6022 (2020); Nat. Mater. **20**, 971 (2021).

